

Introduction to a many-body description of semiconductor lasers

Weng Chow, Sandia National Labs

Outline

- 1) Past: Free-carrier theory
- 2) Present: Many-body description
- 3) Evolving: Quantum-dot gain theory

Collaborations

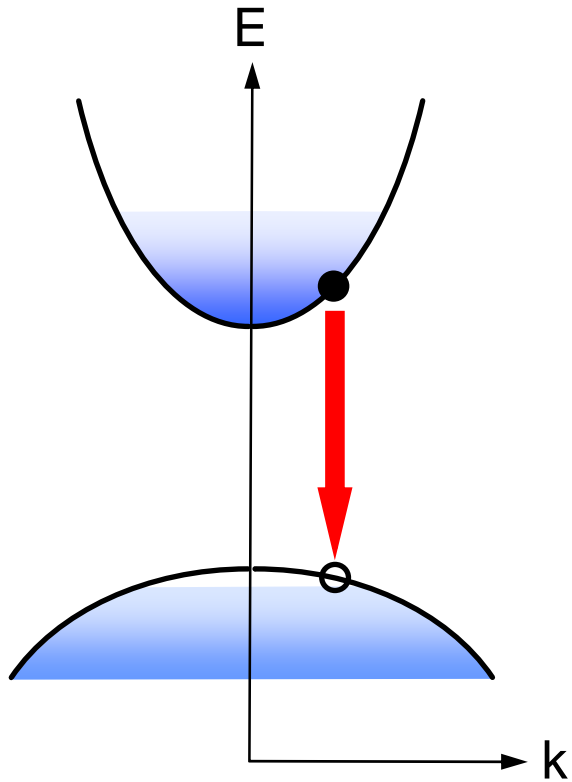
Stephan Koch (Philipps University), Peter Smowton and Peter Blood (Cardiff University), Michael Lorke and Frank Jahnke (Bremen University), Stephan Michael and Christian Schneider (Kaiserslautern University), Andreas Knorr (Technical University, Berlin)

Thanks to:

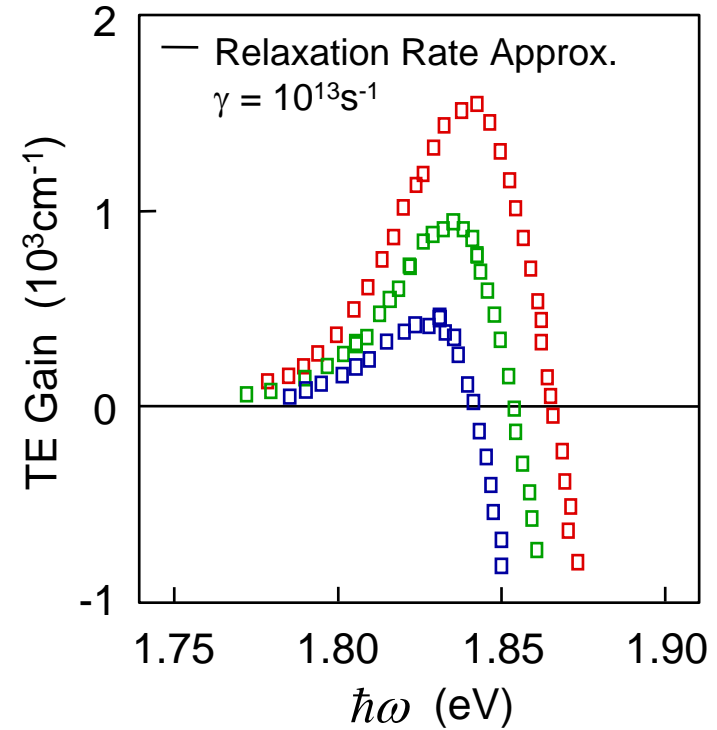
Sandia Laboratory Research and Development (LDRD) program

Humboldt Foundation

Review of semiconductor laser gain



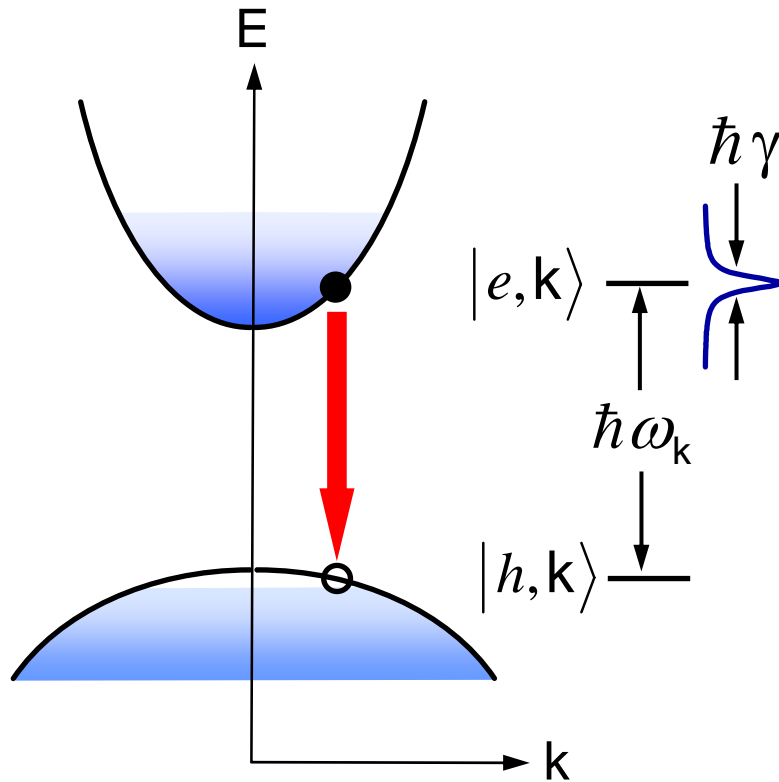
Gain spectrum



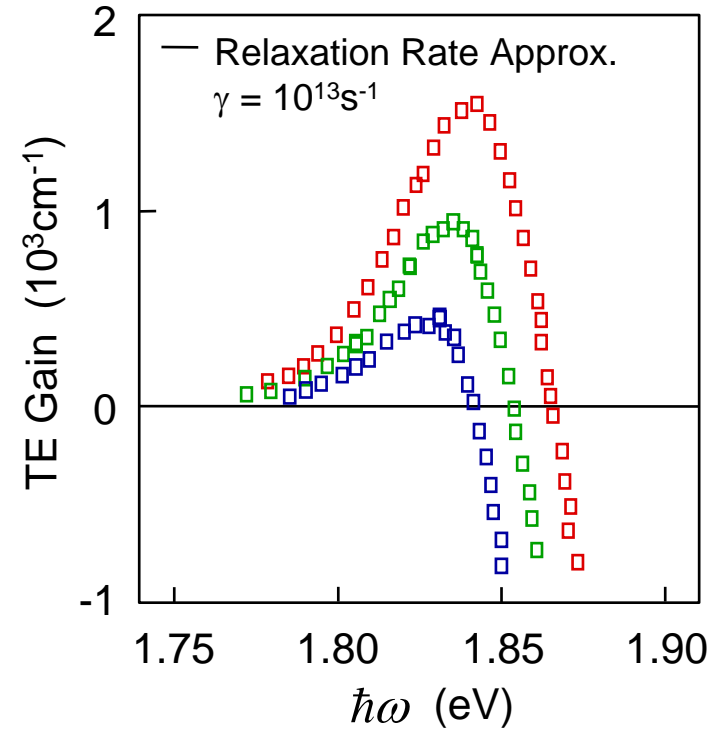
6.8nm $\text{Ga}_{0.41}\text{In}_{0.59}\text{P} / (\text{Al}_{0.5}\text{Ga}_{0.5})_{0.51}\text{In}_{0.49}\text{P}$

Expt. data from P.M. Smowton and P. Blood,
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Free-carrier description

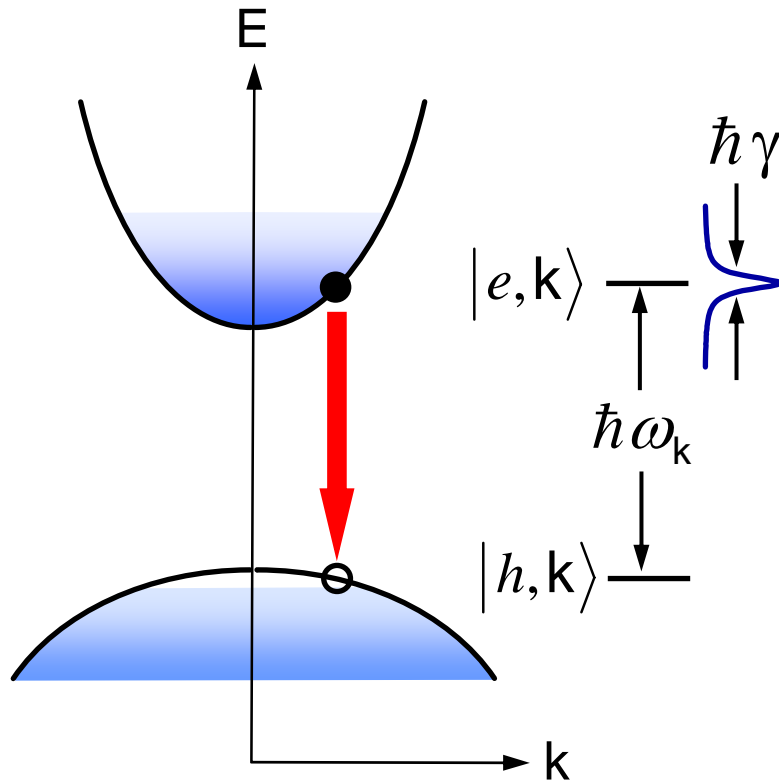
$$\text{Gain} : G \propto \sum_k |\mu_k|^2 \underbrace{[n_{ek} + n_{hk} - 1]}_{\text{Population inversion}} \underbrace{L(\omega_k - \nu)}_{\text{Lineshape function (collision effects)}}$$

Dipole matrix element

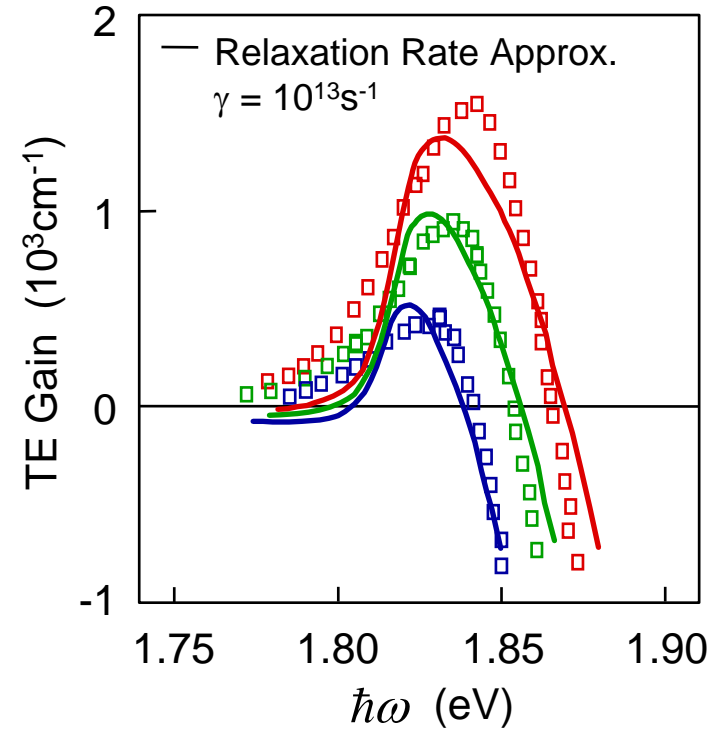
Population inversion

Lineshape function (collision effects)

Review of semiconductor laser gain



Gain spectrum



6.8nm $\text{Ga}_{0.41}\text{In}_{0.59}\text{P} / (\text{Al}_{0.5}\text{Ga}_{0.5})_{0.51}\text{In}_{0.49}\text{P}$

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Free-carrier description

$$\text{Gain} : G \propto \sum_k |\mu_k|^2 \underbrace{[n_{ek} + n_{hk} - 1]}_{\text{Population inversion}} \underbrace{L(\omega_k - \nu)}_{\text{Lineshape function (collision effects)}}$$


Dipole matrix element

Population inversion

Lineshape function (collision effects)

(1) Semiclassical laser theory

$$\frac{d\mathcal{E}}{dz} = g\mathcal{E} = -\frac{\nu}{2\epsilon_0 n c} \text{Im} \left(\sum_k \mu_k \langle b_{-k} a_k \rangle \right)$$

 **Gain**

(2) Equation of motion (Heisenberg Picture)

$$\frac{dO}{dt} = \frac{i}{\hbar} [H, O]$$

(3) Hamiltonian

$$H = \sum_k \left(\varepsilon_{ek} a_k^\dagger a_k + \varepsilon_{hk} b_{-k}^\dagger b_{-k} \right) \quad \text{Kinetic energy}$$
$$- \sum_k \left(\mu_k a_k^\dagger b_{-k}^\dagger + \mu_k^* b_{-k} a_k \right) E(z, t) \quad \text{Light-matter interaction}$$
$$+ \frac{1}{2} \sum_{k, k', q} V_q \left(a_{k+q}^\dagger a_{k'-q}^\dagger a_{k'} a_k + b_{k+q}^\dagger b_{k'-q}^\dagger b_{k'} b_k - 2a_{k+q}^\dagger b_{k'-q}^\dagger b_{k'} a_k \right)$$

Coulomb interaction

(4) Anticommutation relations

$$\left\{ a_k, a_{k'}^\dagger \right\} = \left\{ b_k, b_{k'}^\dagger \right\} = \delta_{k, k'} \quad \left\{ a_k, a_{k'} \right\} = \left\{ b_k, b_{k'} \right\} = 0$$

Quantum statistics: Identical particles

1) Symmetry (or antisymmetry) of wavefunctions

$$\psi(r_1, r_2) = c_1 \phi_a(r_1) \phi_b(r_2) - c_2 \phi_b(r_1) \phi_a(r_2)$$

2) Bookkeeping \longrightarrow Second quantization

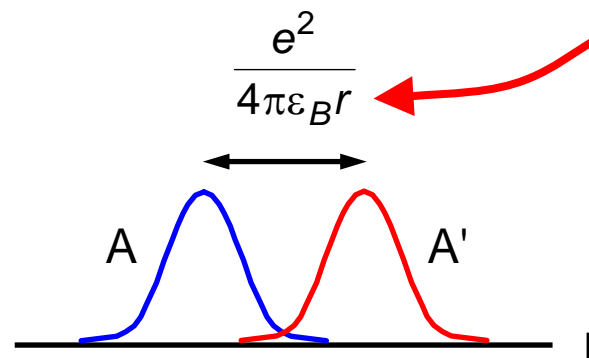
Anticommutation relations

$$\{a_k, a_{k'}^\dagger\} = \{b_k, b_{k'}^\dagger\} = \delta_{k,k'}$$

$$\{a_k, a_{k'}\} = \{b_k, b_{k'}\} = 0$$

$$\{A, B\} = AB + BA$$

3) Quantum statistics and Coulomb interaction



*determines
effective
separation*

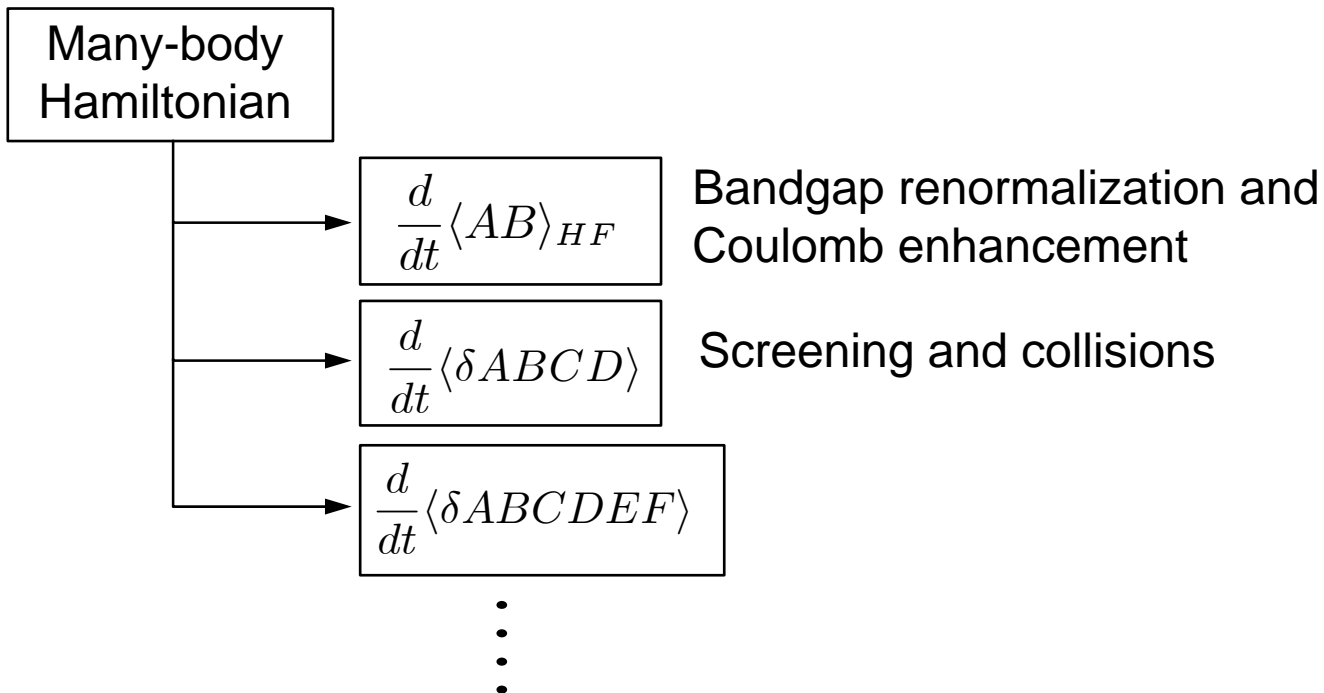
Problem

$$\frac{db_{-k}a_k}{dt} = -i\omega'_k p_k - \frac{i}{\hbar} \mu_k E(z, t) \left[a_k^\dagger a_k + b_{-k}^\dagger b_{-k} - 1 \right] + \frac{i}{\hbar} \sum_{k', q \neq 0} V_q \left[\dots - a_{k'+q}^\dagger b_{-k+q} a_{k'} a_k + \dots \right]$$

Fix

Factorization into products of

$$\begin{cases} p_k = \langle b_{-k} a_k \rangle \\ n_{ek} = \langle a_k^\dagger a_k \rangle \\ n_{hk} = \langle b_{-k}^\dagger b_{-k} \rangle \end{cases}$$



Polarization equation of motion

$$\underbrace{\omega_k^{(0)} - \sum_{k'} V_{kk'} (n_{ek'} + n_{hk'})}_{\text{LHS}} \quad \underbrace{\frac{2\mu_k E}{\hbar} + \sum_{k'} V_{kk'} p_{k'}}_{\text{RHS}}$$
$$\frac{dp_k}{dt} = -i\omega_k p_k - i\Omega_k (n_{ek} + n_{hk} - 1)$$
$$+ S_k^{c-c} + S_k^{c-p}$$

Polarization equation of motion

$$\underbrace{\omega_k^{(0)} - \sum_{k'} V_{kk'} (n_{ek'} + n_{hk'})}_{\text{LHS}} \quad \underbrace{\frac{2\mu_k E}{\hbar} + \sum_{k'} V_{kk'} p_{k'}}_{\text{RHS}}$$

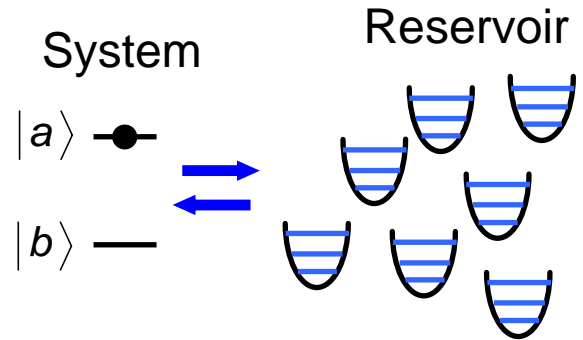
$$\frac{dp_k}{dt} = -i\omega_k p_k - i\Omega_k (n_{ek} + n_{hk} - 1)$$

$$+ S_k^{c-c} + S_k^{c-p}$$

$$- \sigma_k^d p_k$$

Exponential decay

Typical



Polarization equation of motion

$$\underbrace{\omega_k^{(0)} - \sum_{k'} V_{kk'} (n_{ek'} + n_{hk'})}_{\text{Left side}} \quad \underbrace{\frac{2\mu_k E}{\hbar} + \sum_{k'} V_{kk'} p_{k'}}_{\text{Right side}}$$

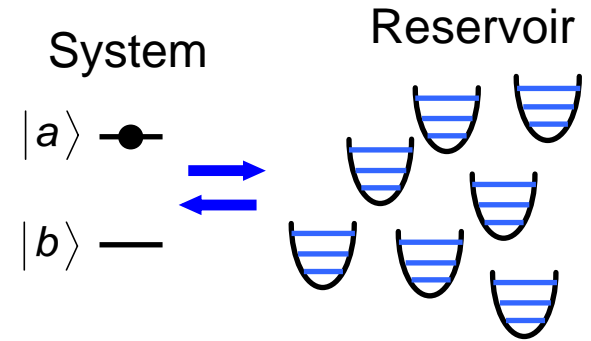
$$\frac{dp_k}{dt} = -i\omega_k p_k - i\Omega_k (n_{ek} + n_{hk} - 1)$$

$$+ S_k^{c-c} + S_k^{c-p}$$

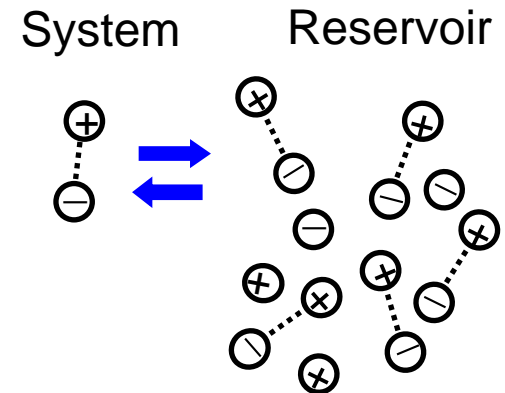
$$-\sigma_k^d p_k + \sum_k \sigma_{k,k'}^{nd} p_{k'}$$

Diagonal **Nondiagonal**

Typical

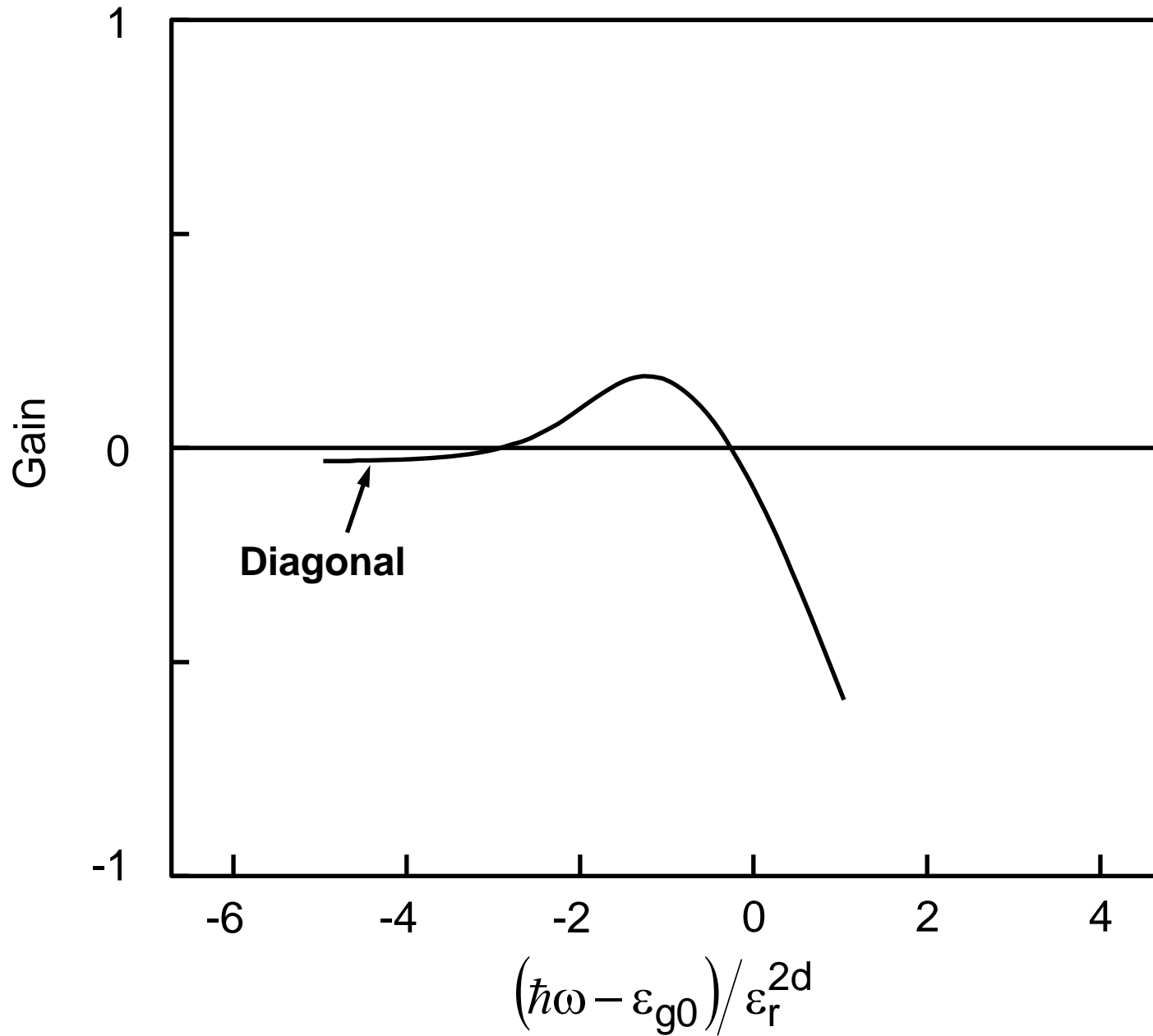


Semiconductor

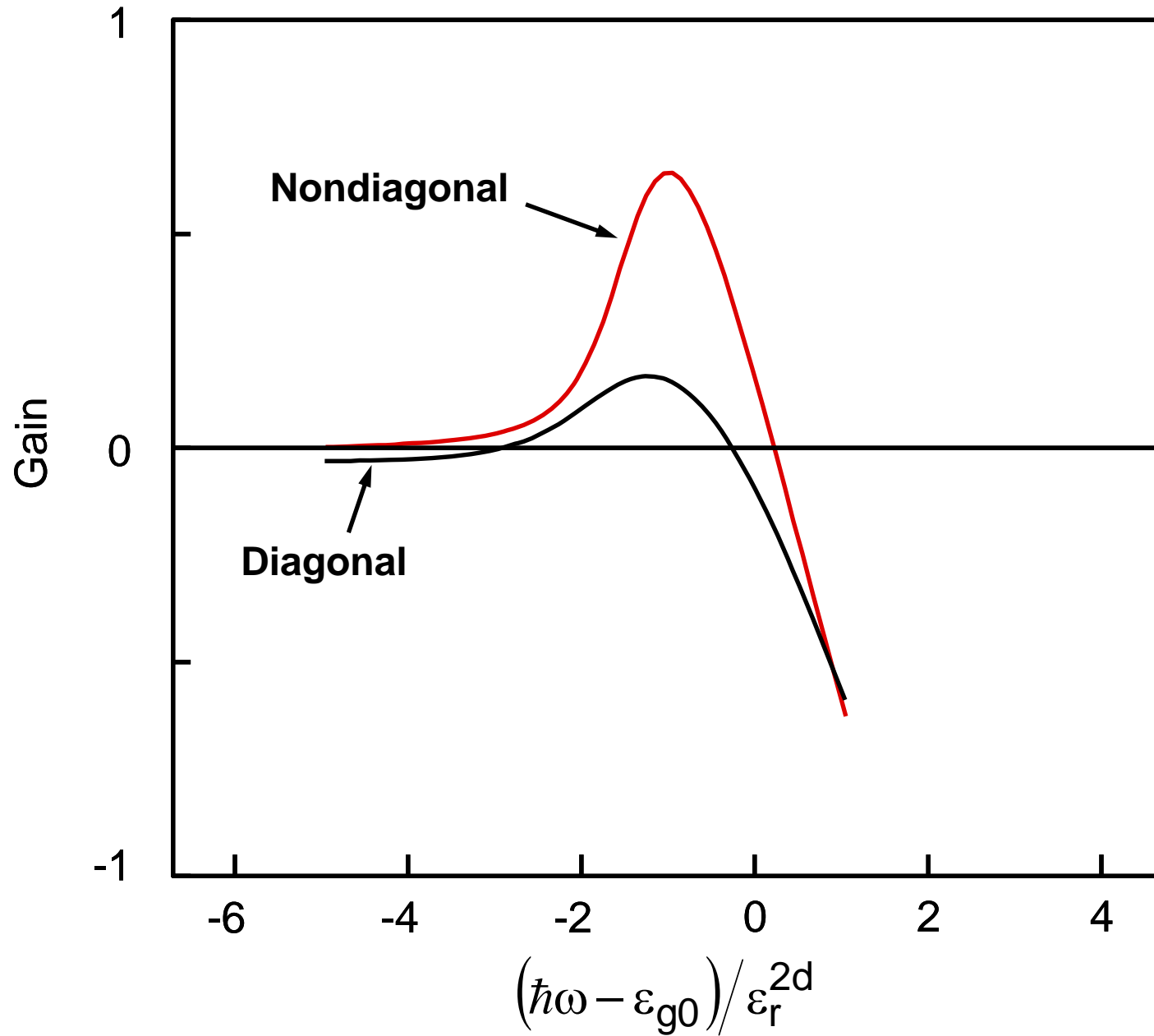


System is its own reservoir

Quantum well gain spectrum

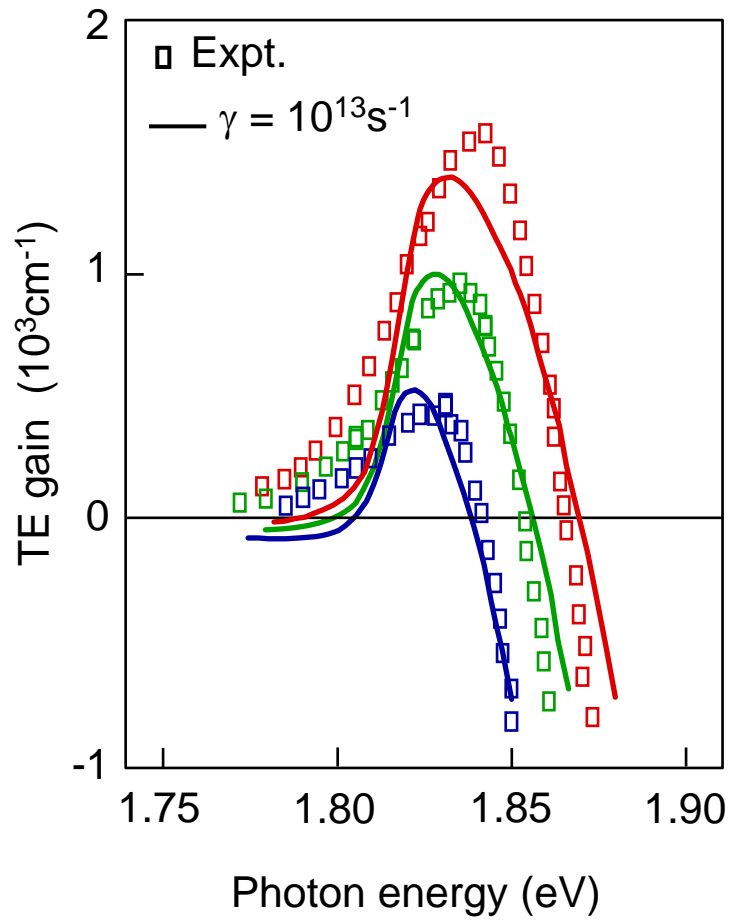


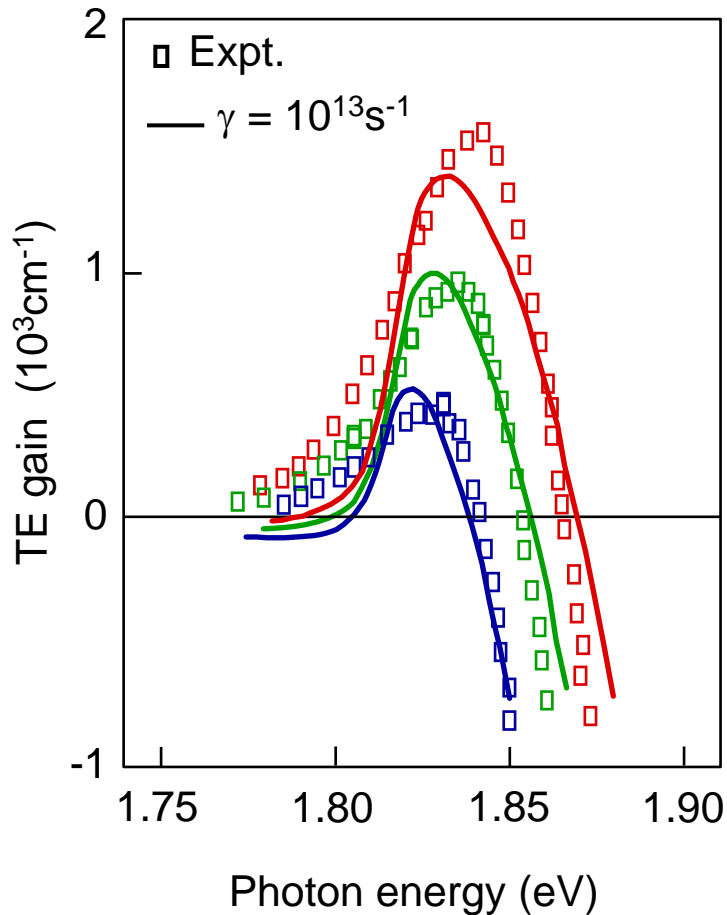
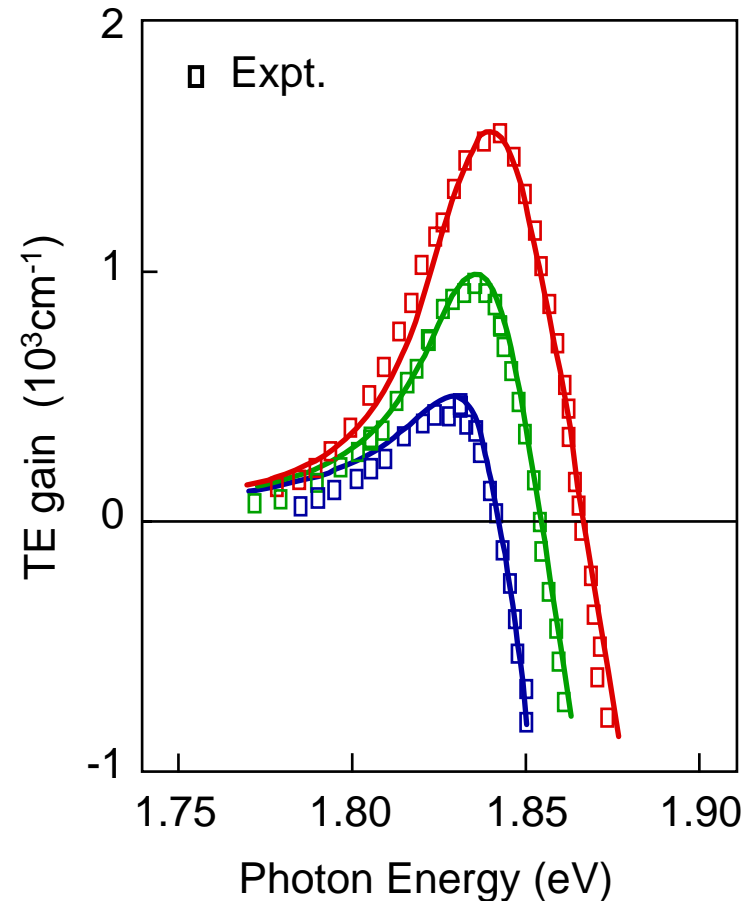
Quantum well gain spectrum



6.8 nm $\text{Ga}_{0.41}\text{In}_{0.59}\text{P} / (\text{Al}_{0.5}\text{Ga}_{0.5})_{0.51}\text{In}_{0.49}\text{P}$

Free-carrier theory



Free-carrier theoryMany-body theory**With many-body effects:**

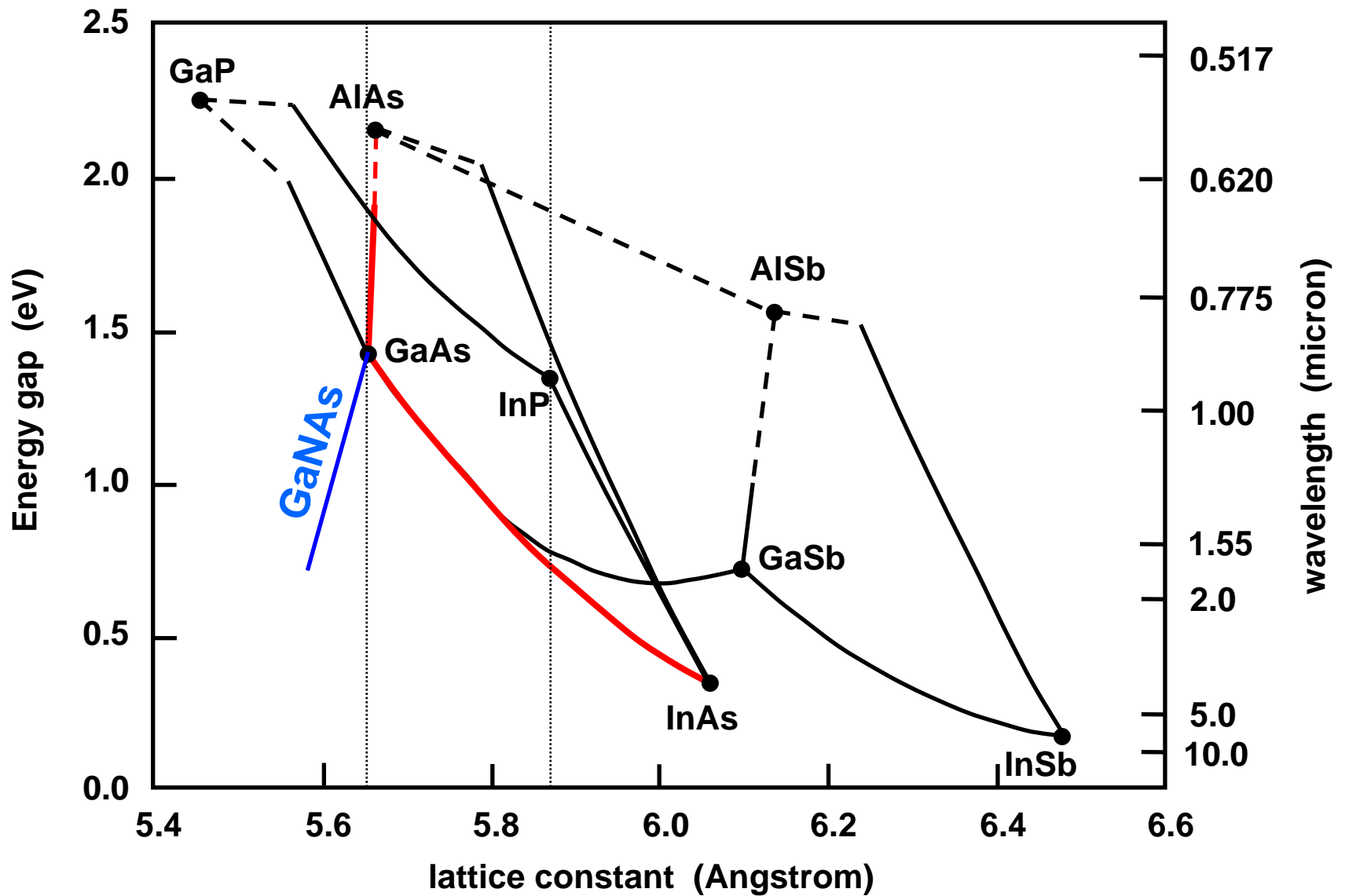
Fix absorption below bandgap

Eliminate dephasing rate as free parameter

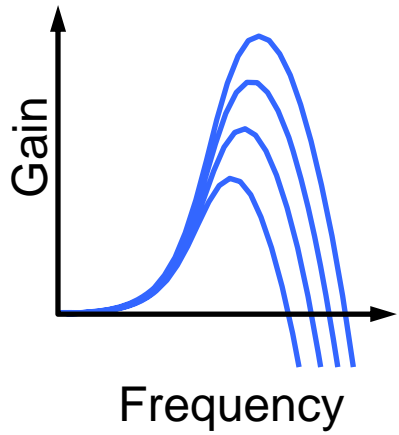
Input: growth sheet and bandstructure parameters

WWC, P. M. Smowton, P. Blood, A. Girdnt, F. Jahnke and S. W. Koch, *APL*.**71**, 157 (1997)

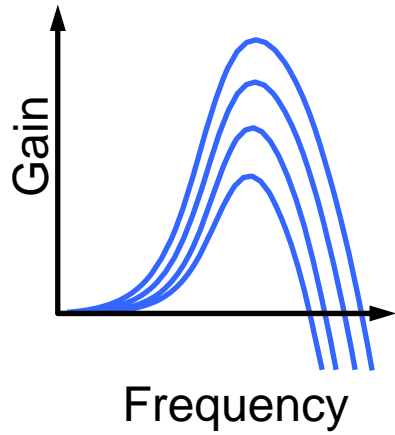
Bandgaps of III-V Alloys (300 K)



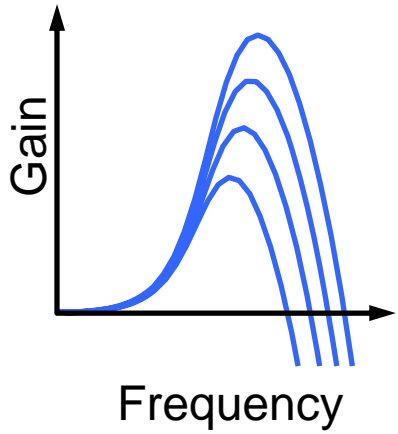
Free-carrier gain



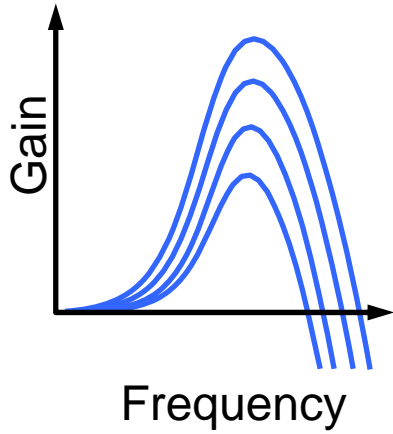
Many-body gain



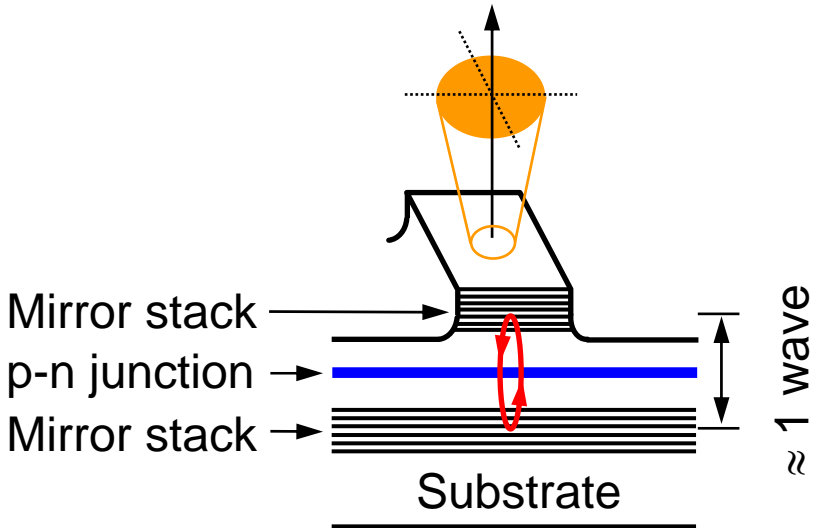
Free-carrier gain



Many-body gain

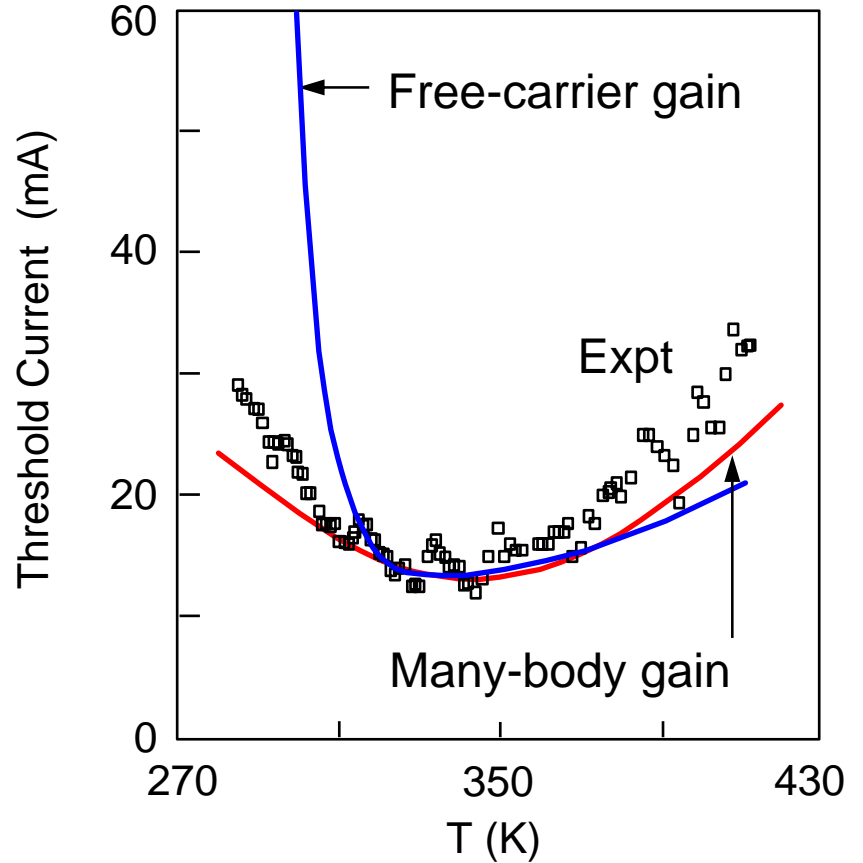


Vertical-Cavity Surface-Emitting Laser (VCSEL)



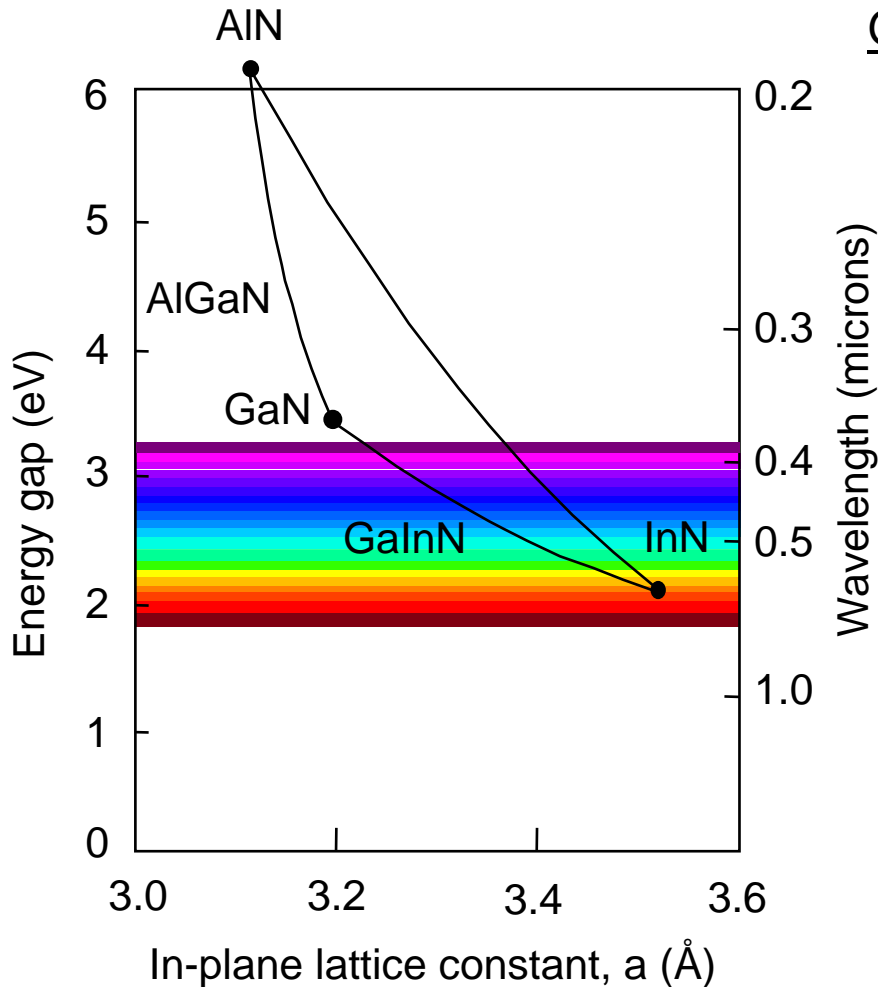
Temperature dependence of VCSEL threshold current

8nm $\text{In}_{0.2}\text{Ga}_{0.8}\text{As}$ -GaAs VCSEL



WWC, Corzine, Young, Coldren
APL 66, 2560 (1995)

Visible and ultraviolet lasers: wide bandgap Group-III Nitrides



Group-III nitride vs. conventional III-V lasers

High exciton binding energy
(23meV for GaN vs. 5meV for GaAs)

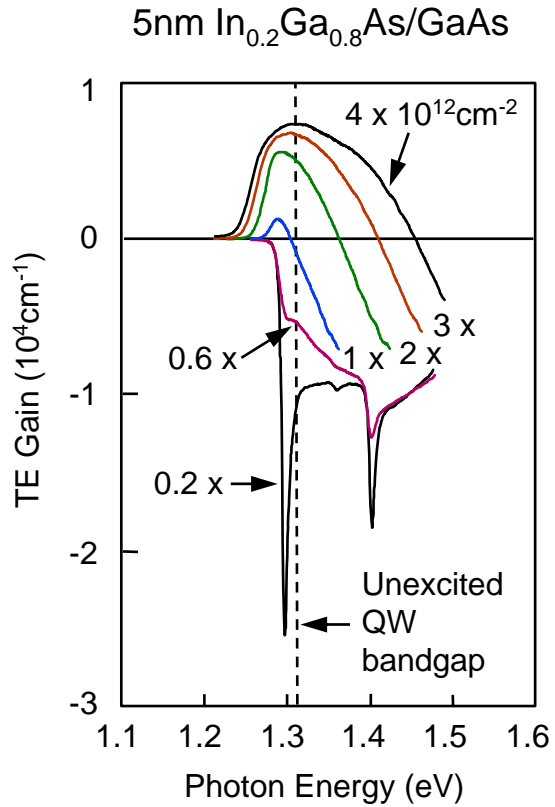
➔ **Strong many-body Coulomb effects**

Wurtzite crystal symmetry

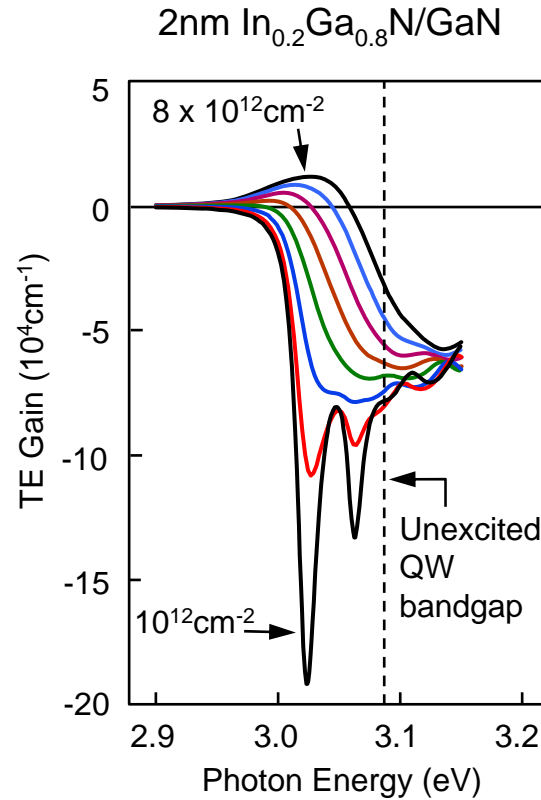
➔ **Strong quantum confined Stark effect**

Strong many-body and bandstructure effects in Wurtzite InGaN

Conventional diode laser

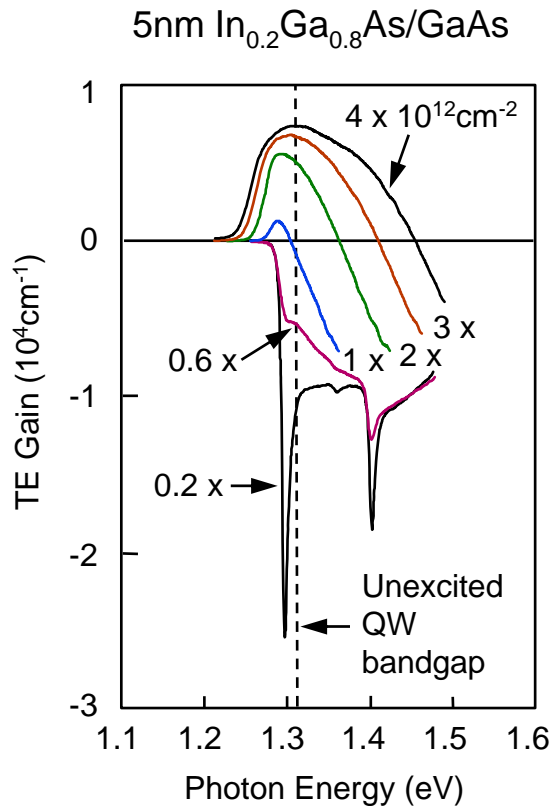


Strong many-body effects

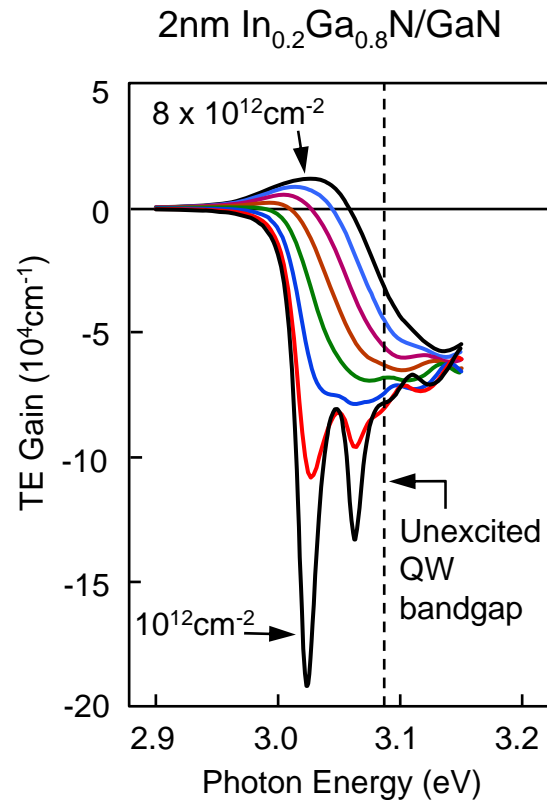


Strong many-body and bandstructure effects in Wurtzite InGaN

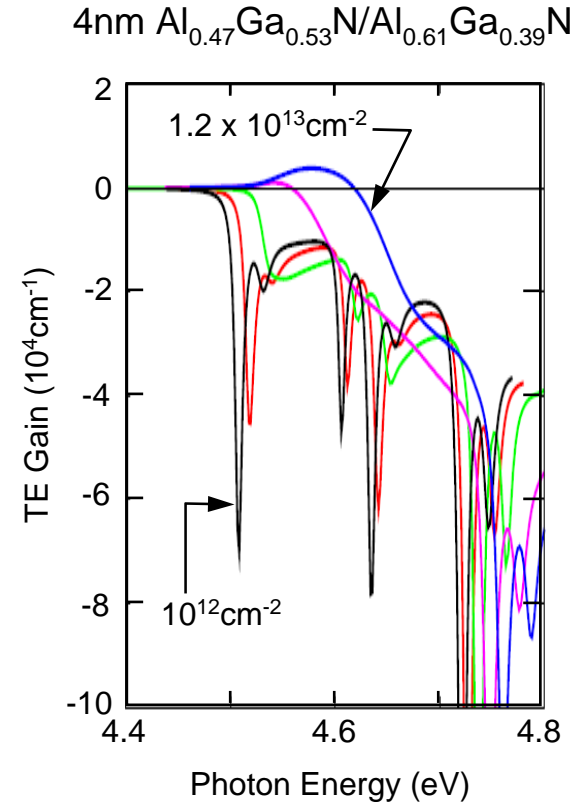
Conventional diode laser



Strong many-body effects



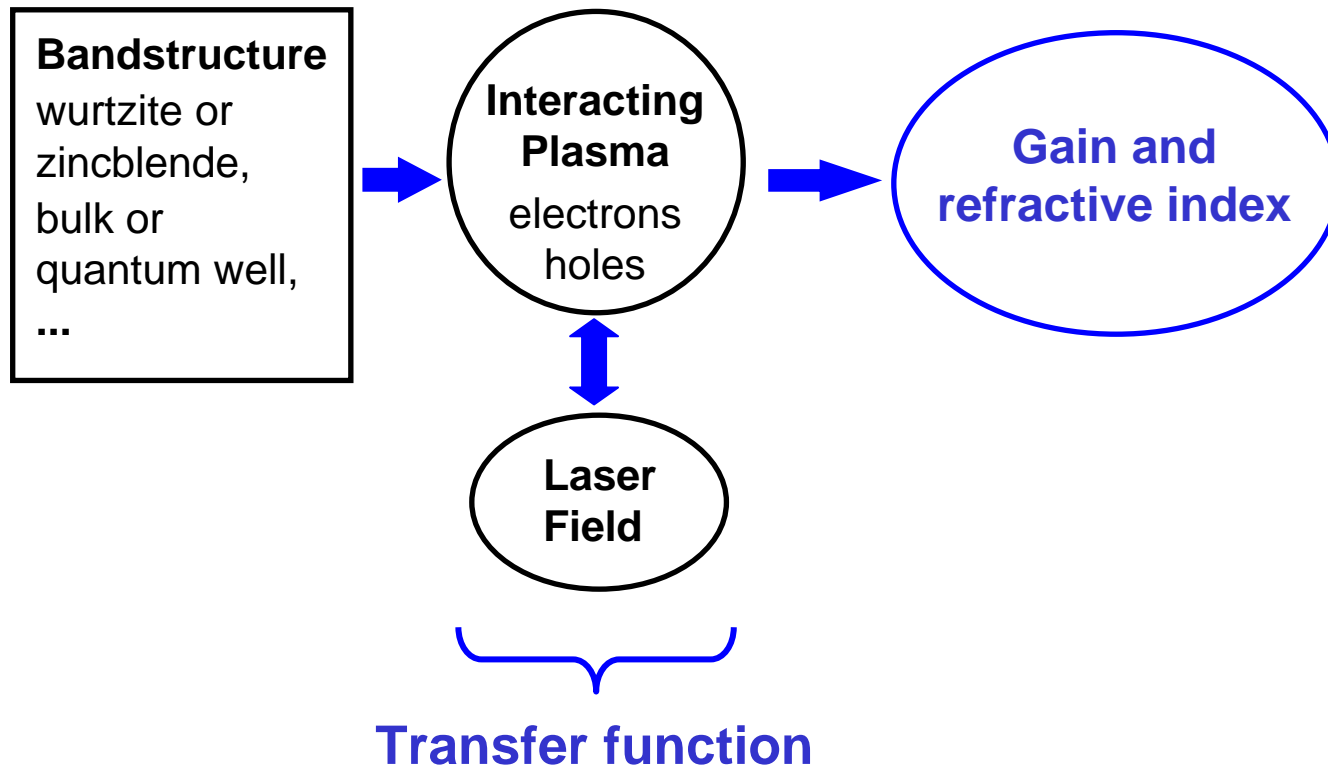
Strong many-body and quantum confined Stark effect



WWC and M. Kneissl (PARC)
 J. Appl. Phys. **98**, 114502 (2005)

T. M. Al Tahtamouni, N. Nepal, J. Y. Lin,
 H. X. Jiang (KSU) and WWC, APL **89**
 131922 (2006)

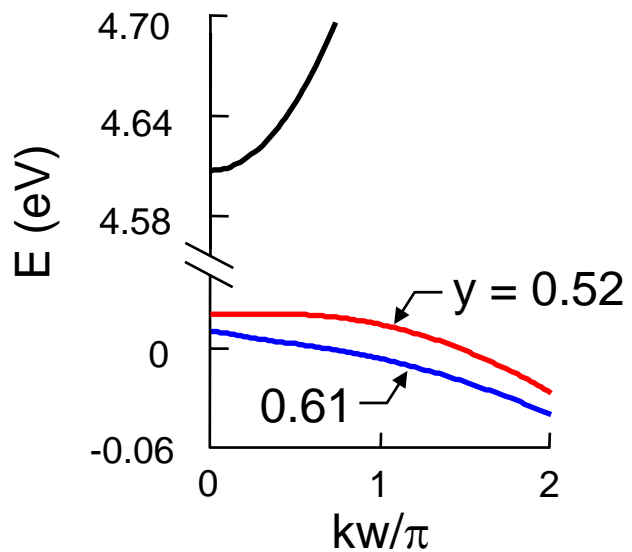
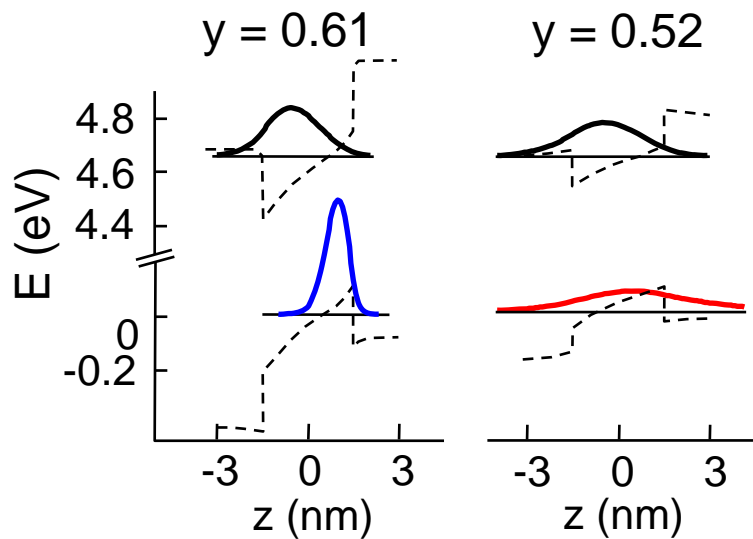
Uses of gain theory



Bandstructure \longrightarrow **Gain properties**

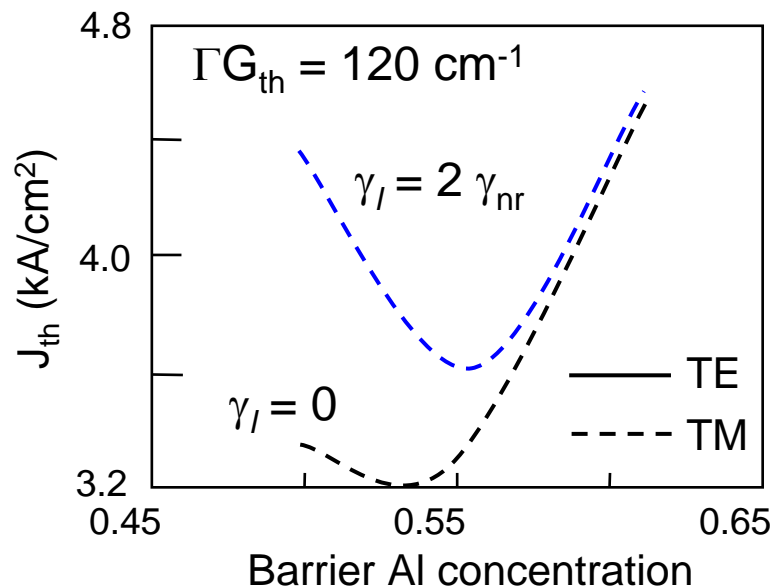
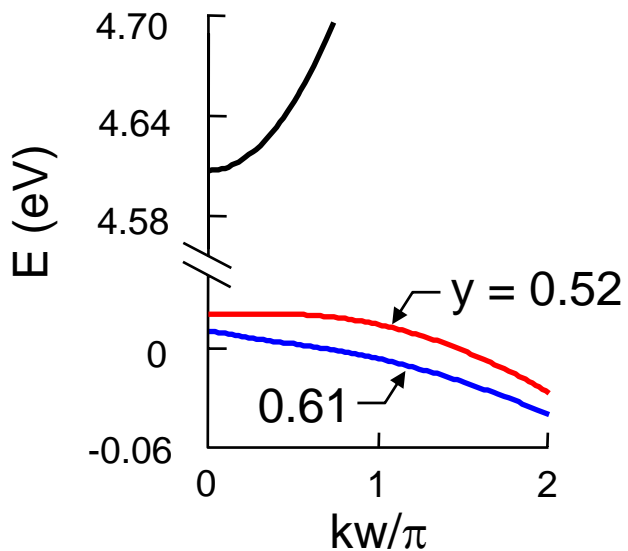
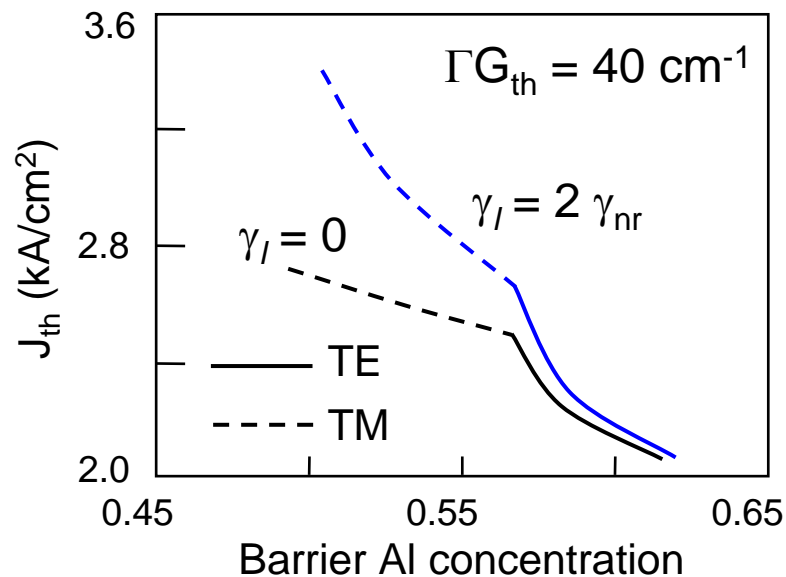
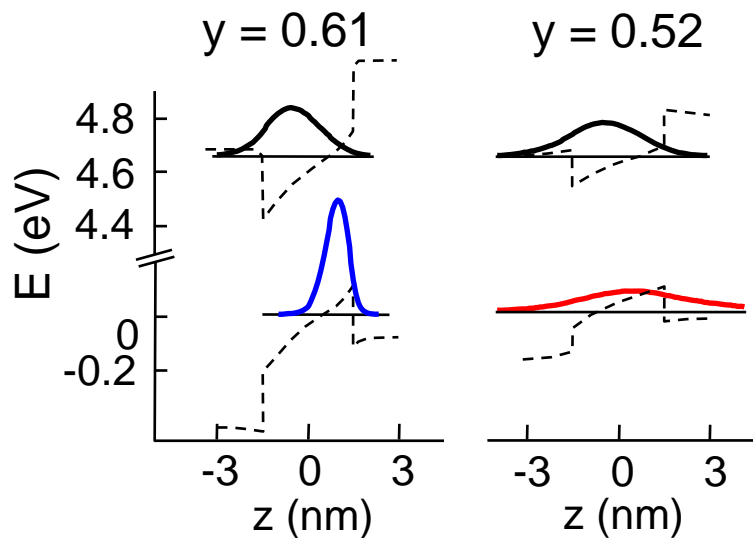
Barrier Influence on Threshold Current Density in AlGaN Lasers

3nm $\text{Al}_{0.46}\text{Ga}_{0.54}\text{N}$ / $\text{Al}_y\text{Ga}_{1-y}\text{N}$, 300K, 280 nm

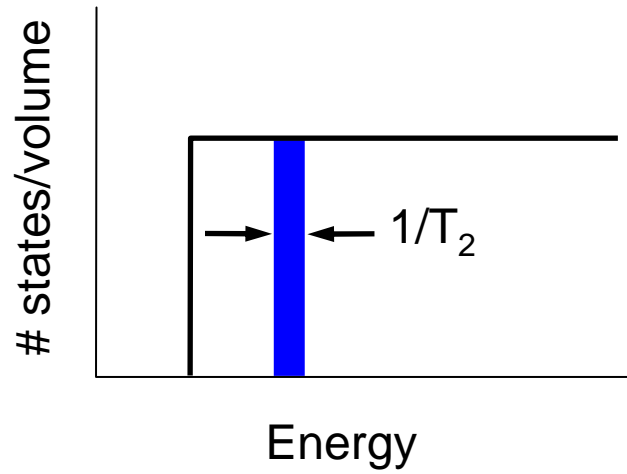
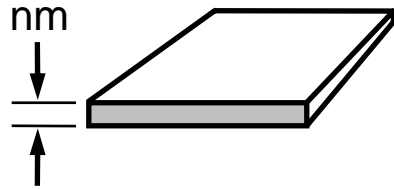


Barrier Influence on Threshold Current Density in AlGaN Lasers

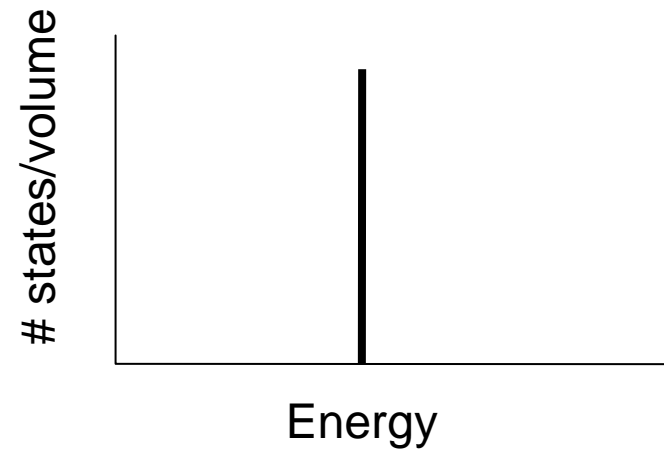
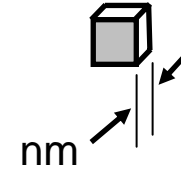
3nm $\text{Al}_{0.46}\text{Ga}_{0.54}\text{N} / \text{Al}_y\text{Ga}_{1-y}\text{N}$, 300K, 280 nm



Quantum well



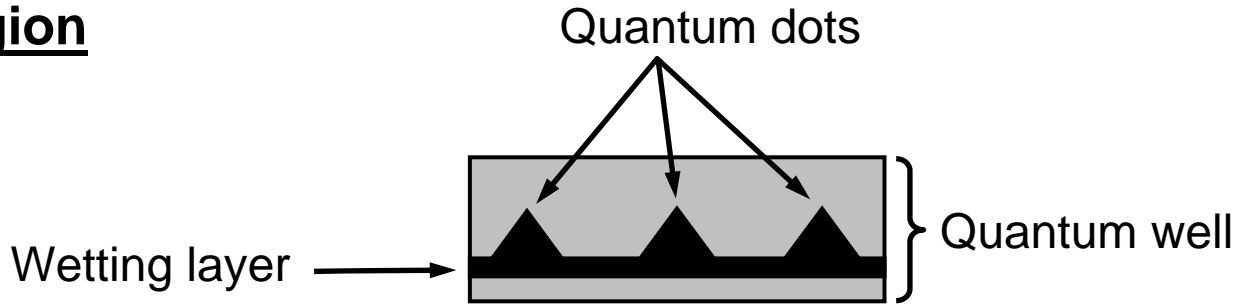
Quantum dot



?

Quantum dot = atom

Active region



Dot-Well Hamiltonian

$$\begin{aligned} H = & \sum_n \varepsilon_n a_n^\dagger a_n + \sum_m \varepsilon_m b_m^\dagger b_m && \text{Single-particle} \\ & - \sum_{n,m} (\mu_{nm} a_n^\dagger b_m^\dagger + \mu_{nm}^* b_m a_n) E(z, t) && \text{Light-matter interaction} \\ & + \frac{1}{2} \sum_{n,m,r,s} W_{nm}^{rs} a_r^\dagger a_s^\dagger a_m a_n + \frac{1}{2} \sum_{n,m,r,s} W_{nm}^{rs} b_r^\dagger b_s^\dagger b_m b_n - \sum_{n,m,r,s} W_{nm}^{rs} a_r^\dagger b_s^\dagger b_m a_n \\ & + \sum_{n,q} \hbar G_q a_{n+q}^\dagger a_n (b_q + b_{-q}^\dagger) && \text{Carrier - phonon} \quad \text{Carrier - carrier} \end{aligned}$$

Coulomb matrix element

$$W_{nk}^{rk'} = \int d^3r d^3r' \phi_r^*(\mathbf{r}) \phi_{k'}^*(\mathbf{r}') \frac{e^2}{4\pi\varepsilon_b |\mathbf{r} - \mathbf{r}'|} \phi_k(\mathbf{r}') \phi_n(\mathbf{r})$$

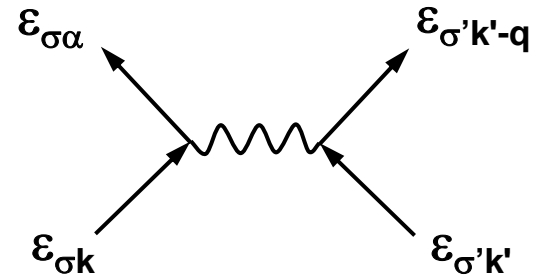
Complications with quantum dots

$$\frac{dp_\alpha}{dt} = -i\omega_\alpha p_\alpha - i\Omega_\alpha (n_{e\alpha} + n_{h\alpha} - 1)$$

$$+ S_\alpha^{c-c} + S_\alpha^{c-p}$$



Carrier-carrier scattering



(diagonal contribution)

$$-p_\alpha \frac{2}{\hbar} \sum_{\mathbf{k}'} \sum_{\mathbf{q}} \sum_{\sigma} \sum_{\sigma'} \sum_{\mathbf{k}} \xi_{\alpha k} \left[W_q^2 - \frac{\delta_{\sigma, \sigma'}}{2} W_q W_{|\mathbf{k}' - \mathbf{q} - \mathbf{k}|} \right] \frac{1}{\Delta - i\delta\varepsilon}$$

$$\times \left[n_{|\mathbf{k}' - \mathbf{q}| \sigma'} (1 - n_{k' \sigma'}) (1 - n_{k \sigma}) + (1 - n_{|\mathbf{k}' - \mathbf{q}| \sigma'}) n_{k' \sigma'} n_{k \sigma} \right]$$

$$\delta\varepsilon = \hbar\omega - \underbrace{\varepsilon_{\bar{\sigma}\alpha} - \varepsilon_{\sigma k} + \varepsilon_{\sigma' k' - q} - \varepsilon_{\sigma' k'}}_{\text{Renormalized energies}}$$

Also, memory effects

Renormalized energies

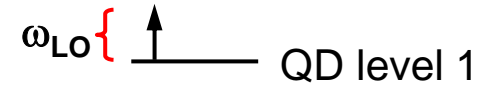
Complications with quantum dots

$$\frac{dp_\alpha}{dt} = -i\omega_\alpha p_\alpha - i\Omega_\alpha (n_{e\alpha} + n_{h\alpha} - 1)$$

$$+ S_\alpha^{c-c} + S_\alpha^{c-p}$$



Carrier-phonon scattering



Complications with quantum dots

$$\frac{dp_\alpha}{dt} = -i\omega_\alpha p_\alpha - i\Omega_\alpha (n_{e\alpha} + n_{h\alpha} - 1)$$

$$+ S_\alpha^{c-c} + S_\alpha^{c-p}$$

—— QD level 2

ω_{LO} { \uparrow — QD level 1

Carrier-phonon scattering

Nonperturbative treatment

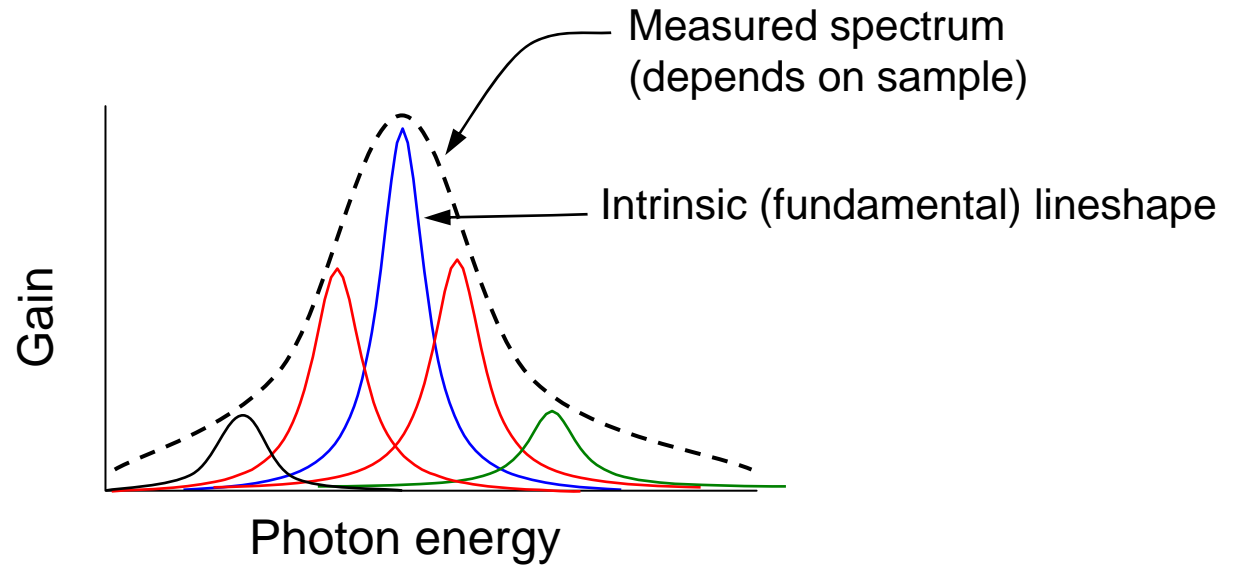
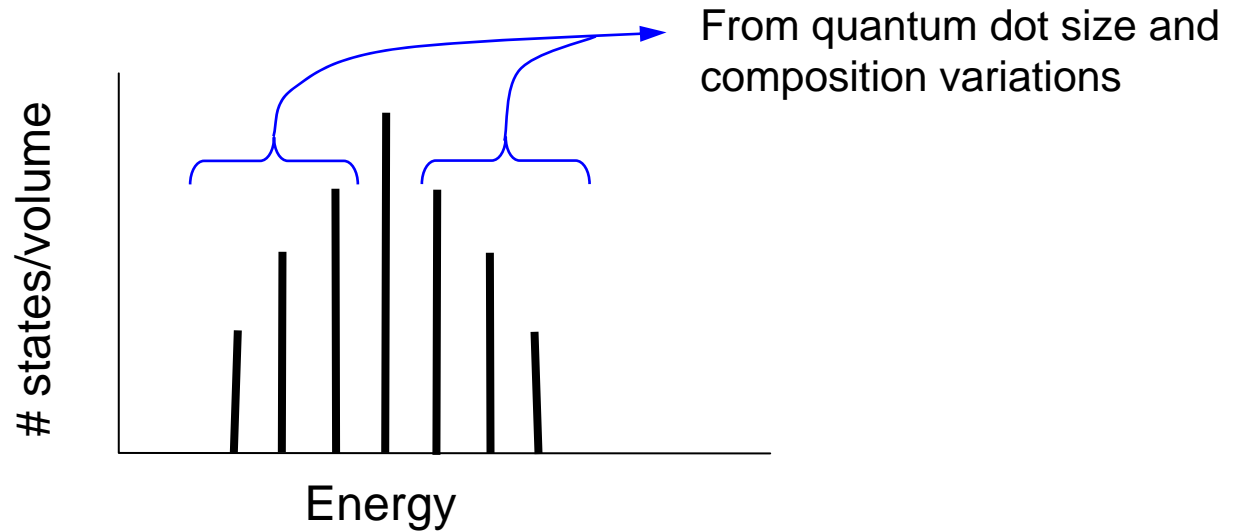
$$-p_\alpha \int_0^\infty dt' \sum_{\mathbf{k}} \left| \frac{M_{\alpha\mathbf{k}}}{\hbar} \right|^2 \sum_{\sigma=e,h} G_{k\sigma}(t') G_{\alpha\sigma}(t') e^{i(\varepsilon_{\alpha\sigma} - \varepsilon_{k\sigma})t'/\hbar}$$

$$\times \left\{ (1 - n_{k\sigma}) \left[n_{LO} e^{i\omega_{LO}t'} + (n_{LO} + 1) e^{-i\omega_{LO}t'} \right] + n_{k\sigma} \left[(n_{LO} + 1) e^{i\omega_{LO}t'} + n_{LO} e^{-i\omega_{LO}t'} \right] \right\}$$

(diagonal contribution)

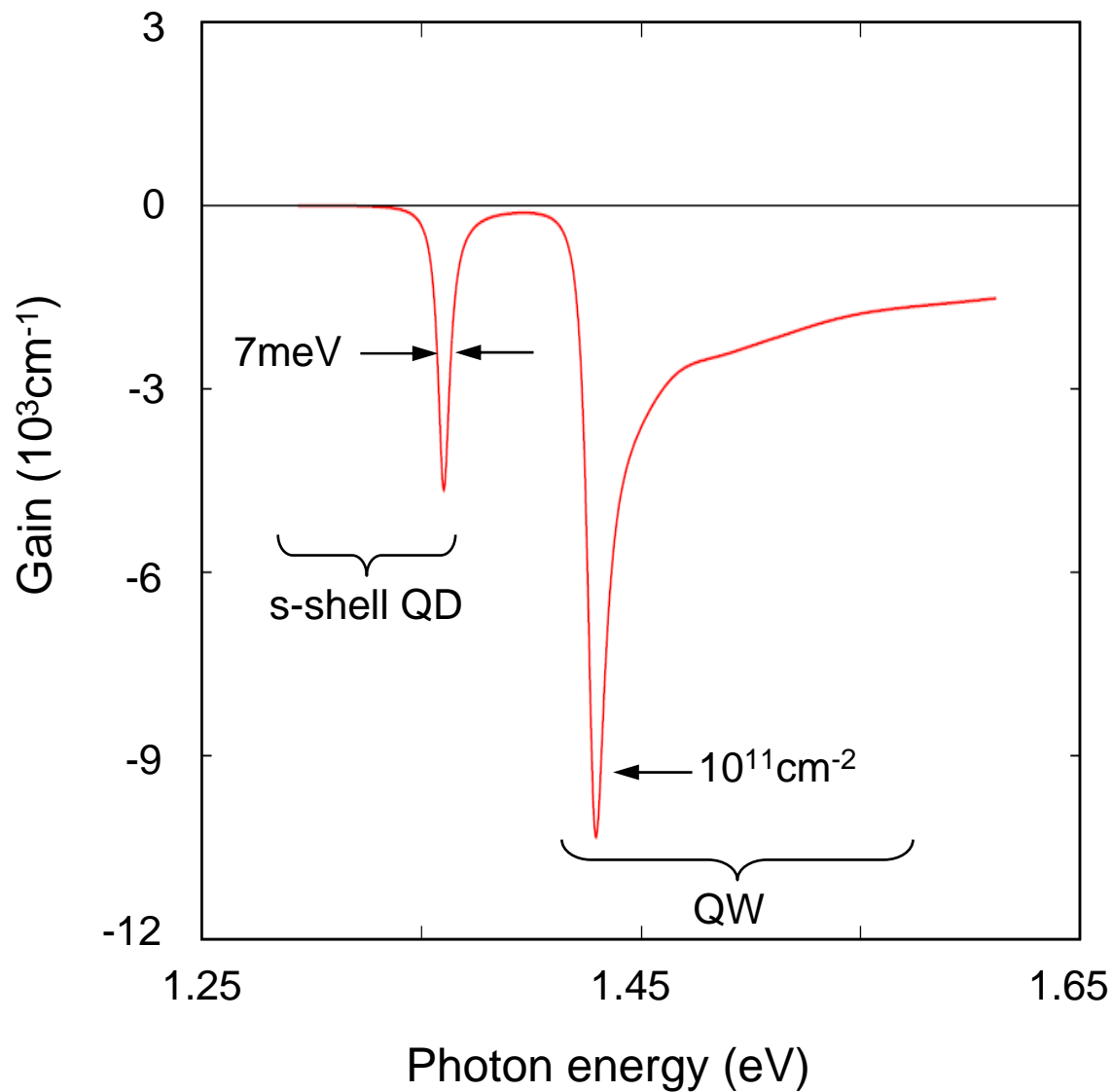
$$\frac{dG_{\alpha\sigma}(t)}{dt} = - \int_0^t dt' \left[(n_{LO} + 1) e^{-i\omega_{LO}t'} + n_{LO} e^{i\omega_{LO}t'} \right] \\ \times G_{\alpha\sigma}(t - t') \sum_{\beta\sigma'} \frac{|M_{\alpha\sigma, \beta\sigma'}|^2}{\hbar^2} G_{\beta\sigma'}(t') \exp \left[i \frac{\varepsilon_{\alpha\sigma} - \varepsilon_{\beta\sigma'}}{\hbar} t' \right]$$

Inhomogeneous broadening



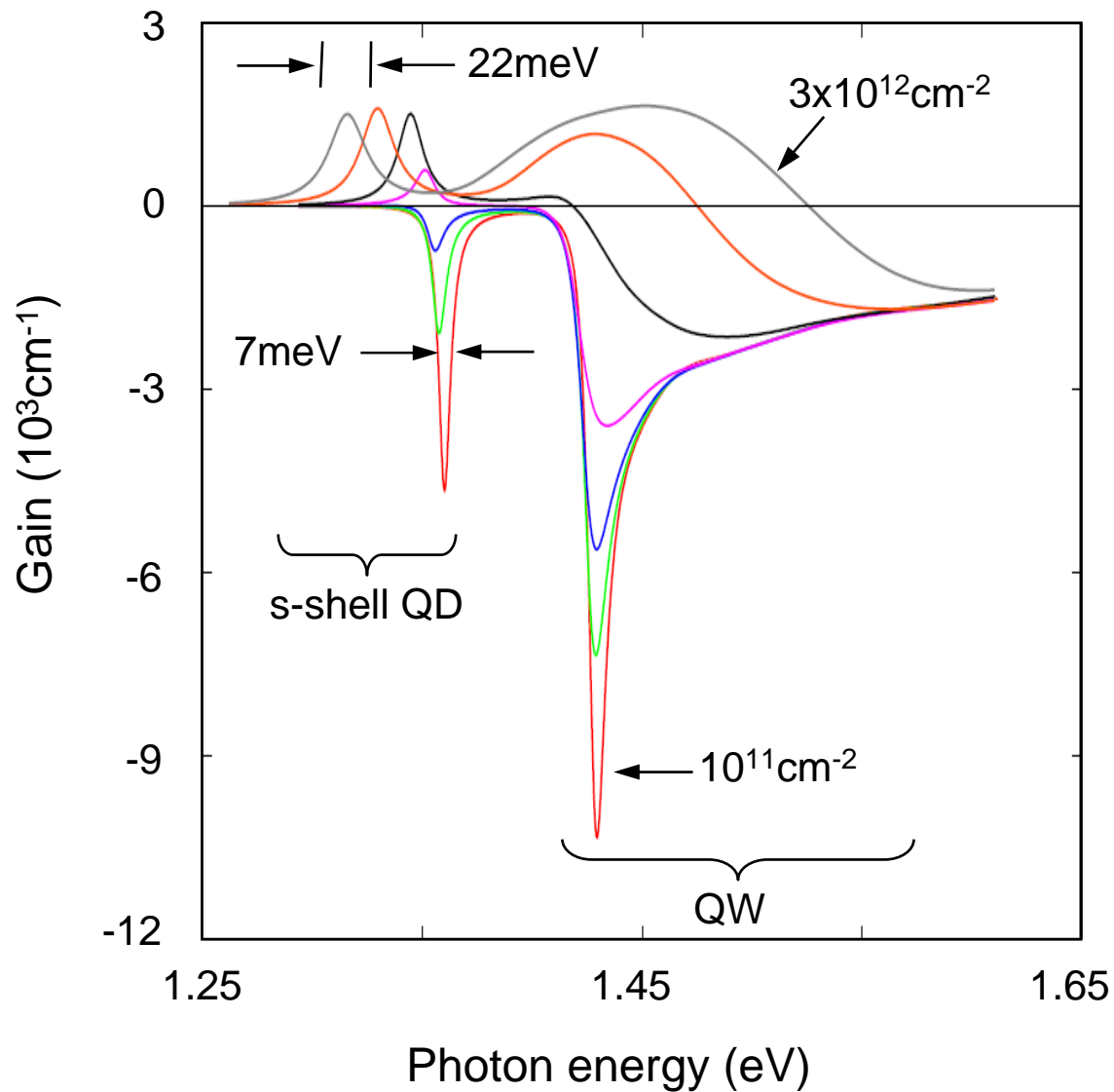
Homogeneously (intrinsically) broadened quantum-dot absorption and gain

InGaAs QD, $N_D=5 \times 10^{10} \text{cm}^{-2}$, $T=300\text{K}$

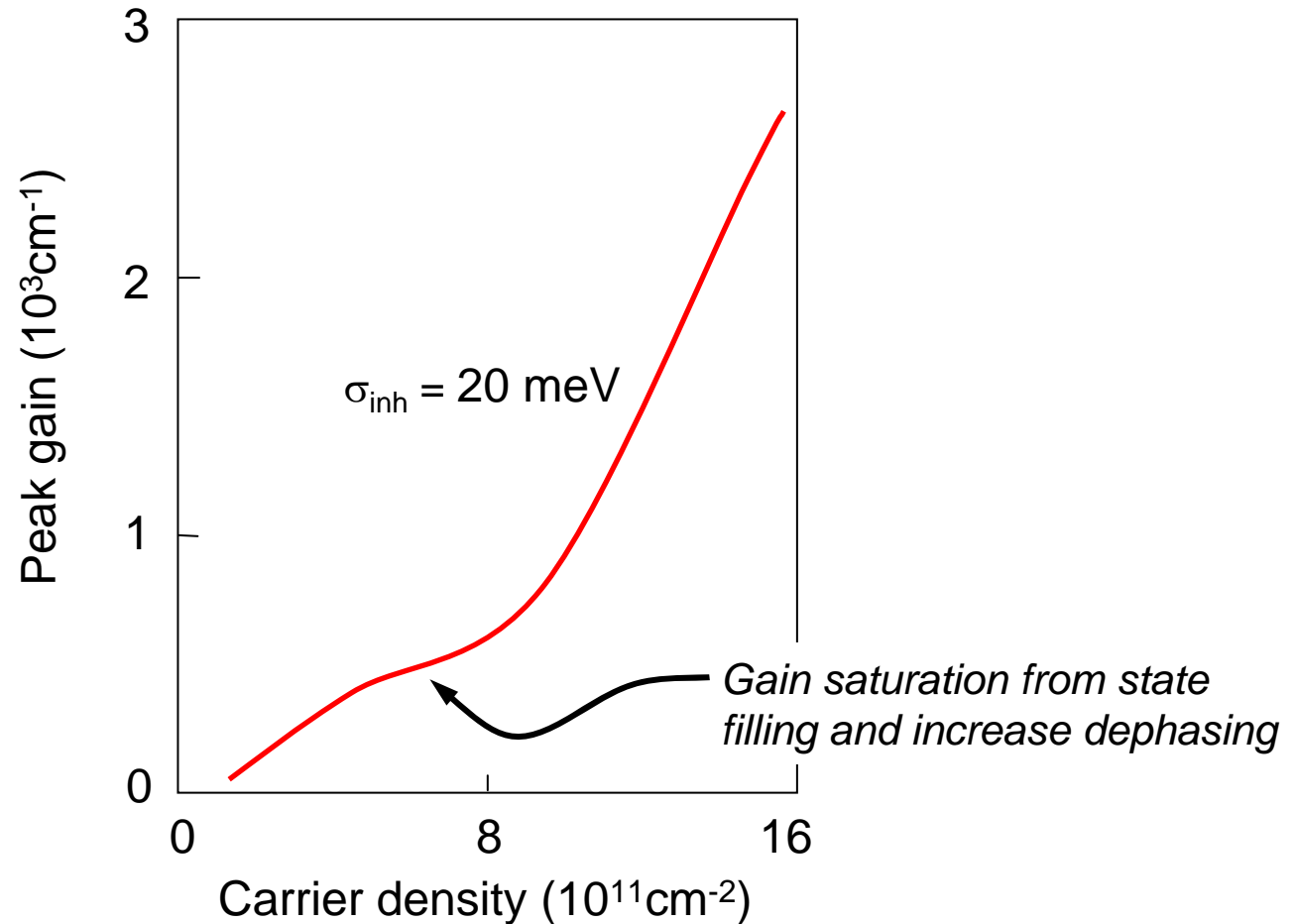


Homogeneously (intrinsically) broadened quantum-dot absorption and gain

InGaAs QD, $N_D=5 \times 10^{10} \text{cm}^{-2}$, $T=300\text{K}$



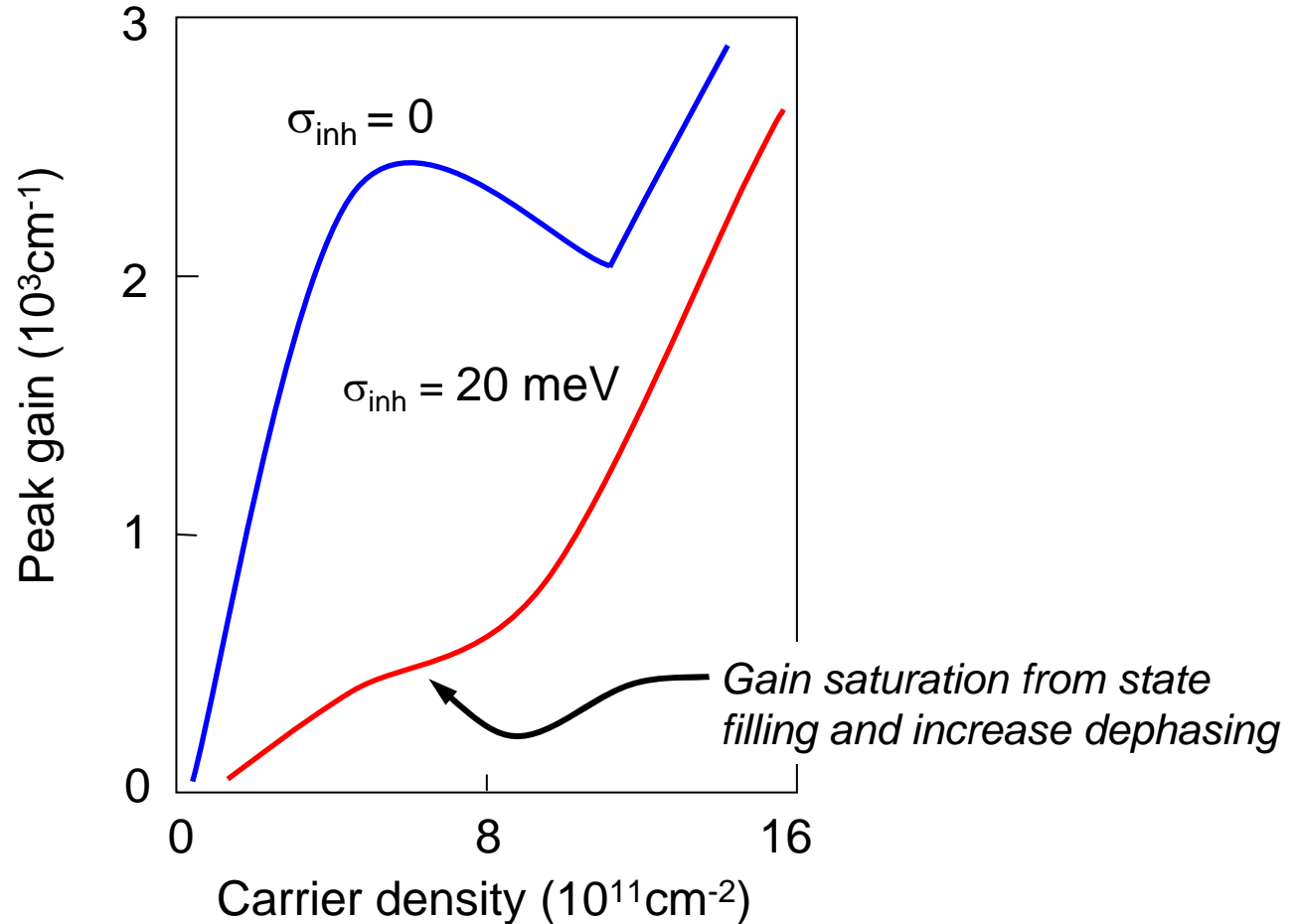
Quantum-dot gain saturation



Lorke, WWC, Nielsen, Seebeck, Gartner and Jahnke, PRB **74**, 035334 (2006)

Lorke, Jahnke, WWC, APL **90**, 051112 (2007)

Quantum-dot gain saturation

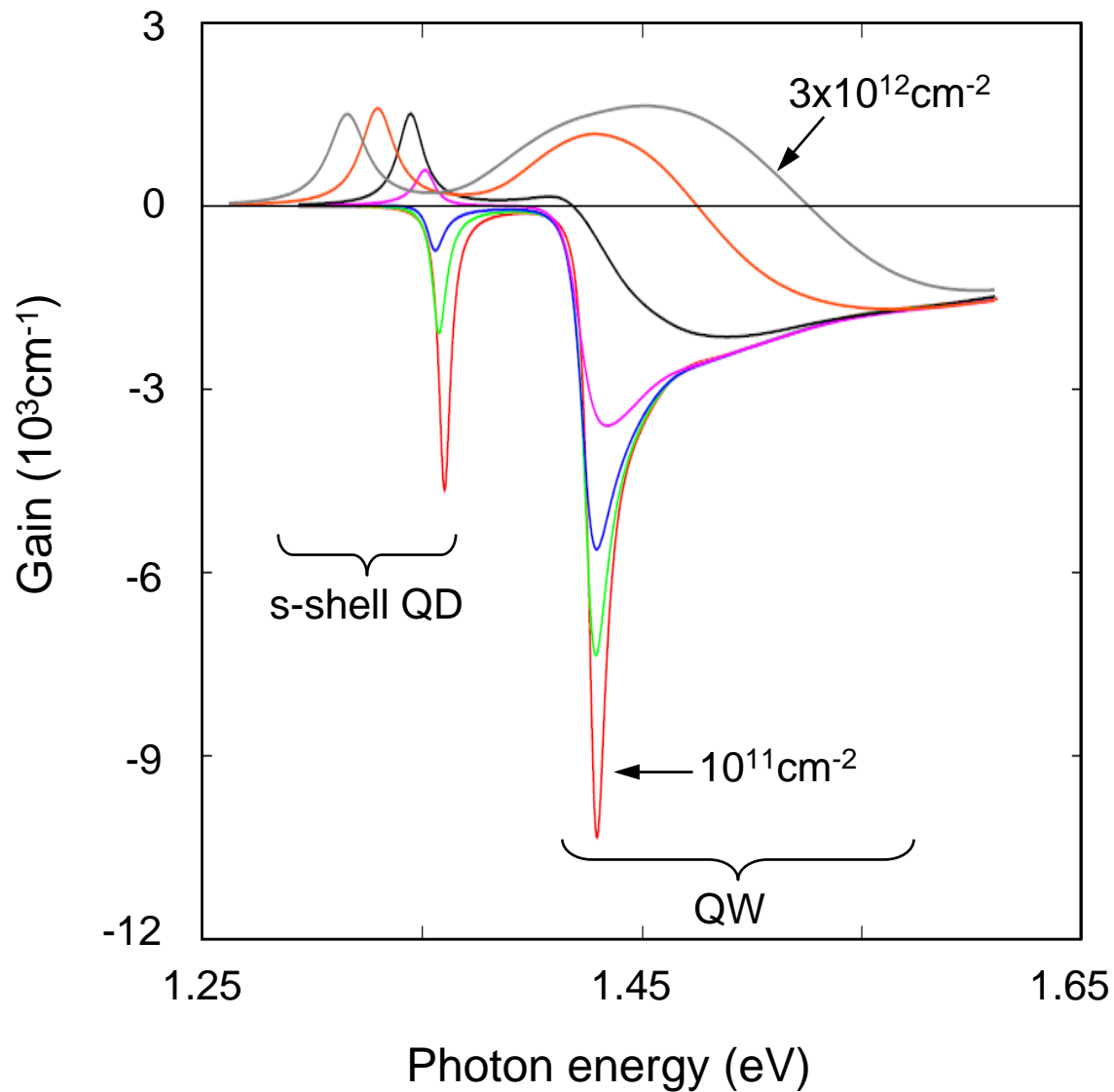


Lorke, WWC, Nielsen, Seebeck, Gartner and Jahnke, PRB **74**, 035334 (2006)

Lorke, Jahnke, WWC, APL **90**, 051112 (2007)

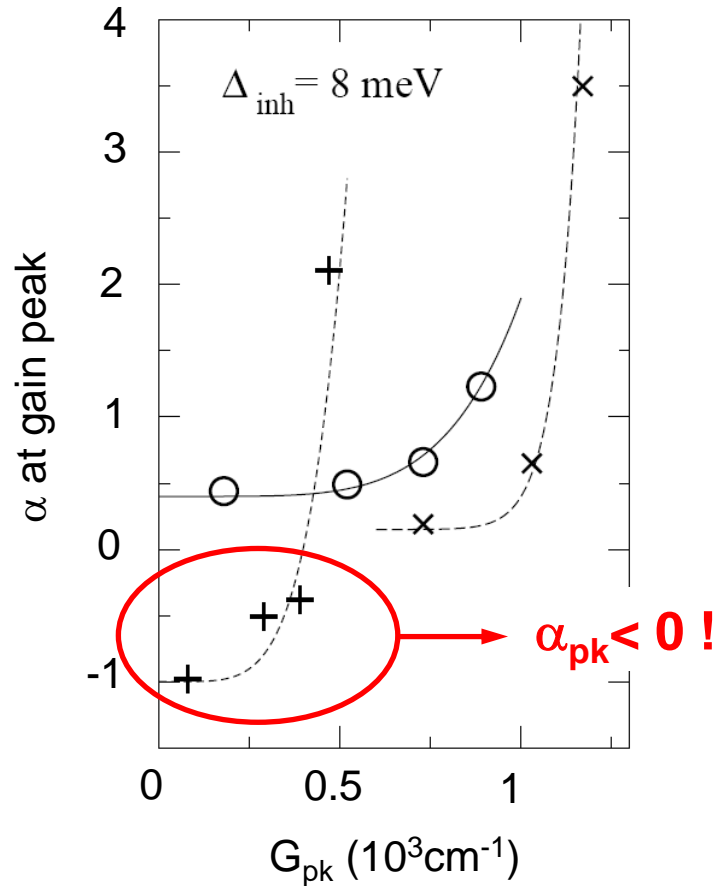
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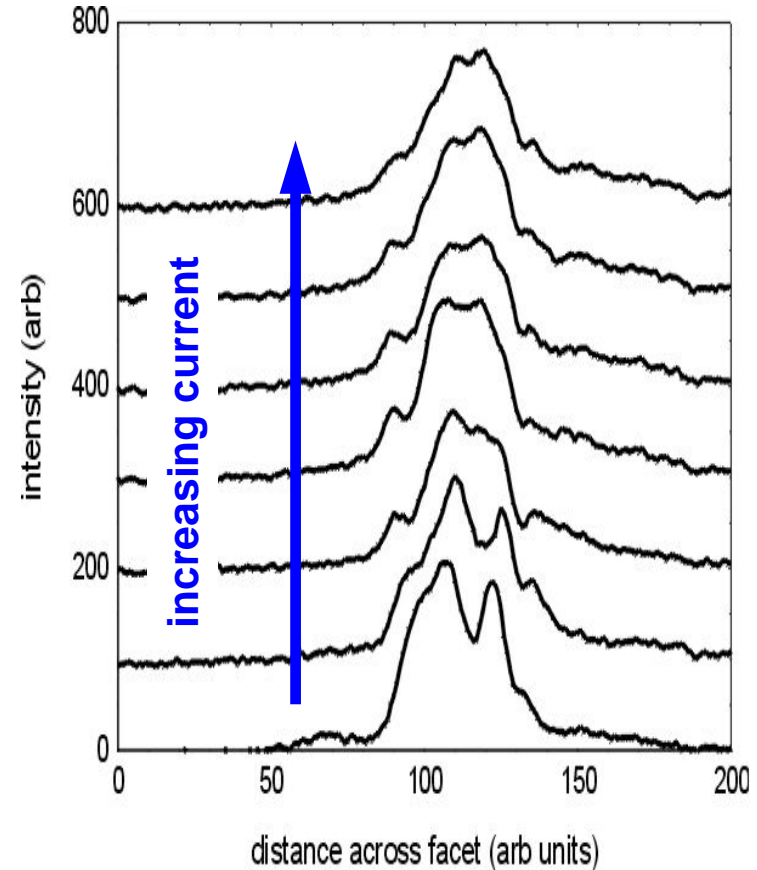
Linewidth enhancement (antiguinding) factor in quantum dots

Many-body calculation



Near field measurements

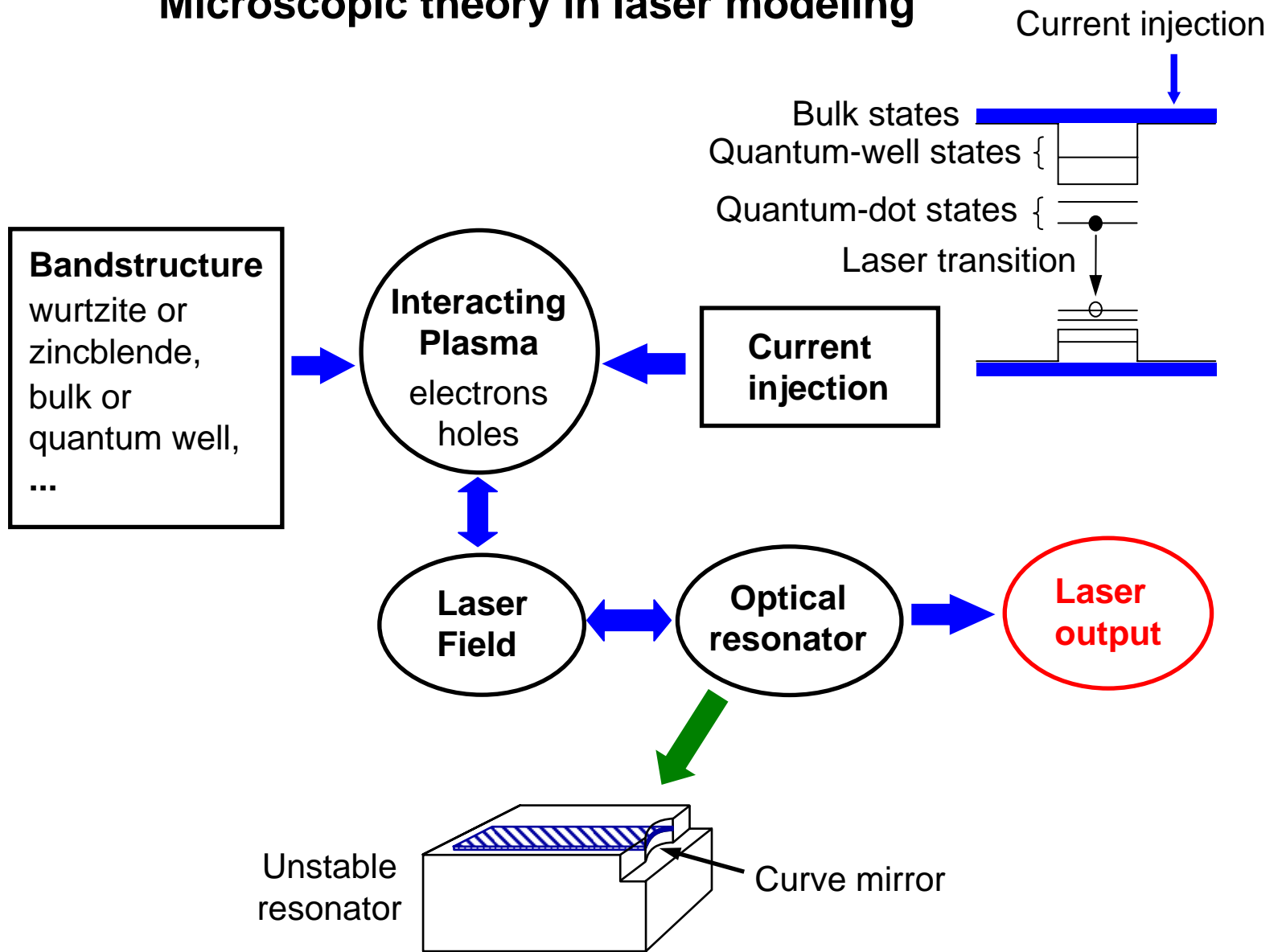
P. M. Snowton, et al, APL 81, 3251, 2002



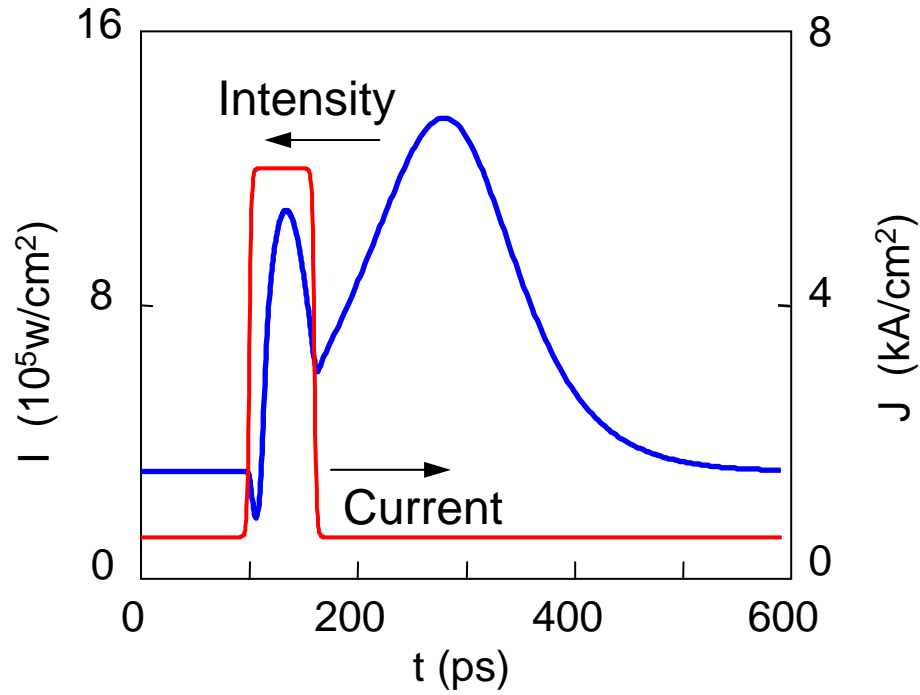
Lorke, Jahnke, WWC, APL **90**, 051112 (2007)

Schneider, WWC, Koch, PRB **66**, Rapid Commun. 41310, 2002

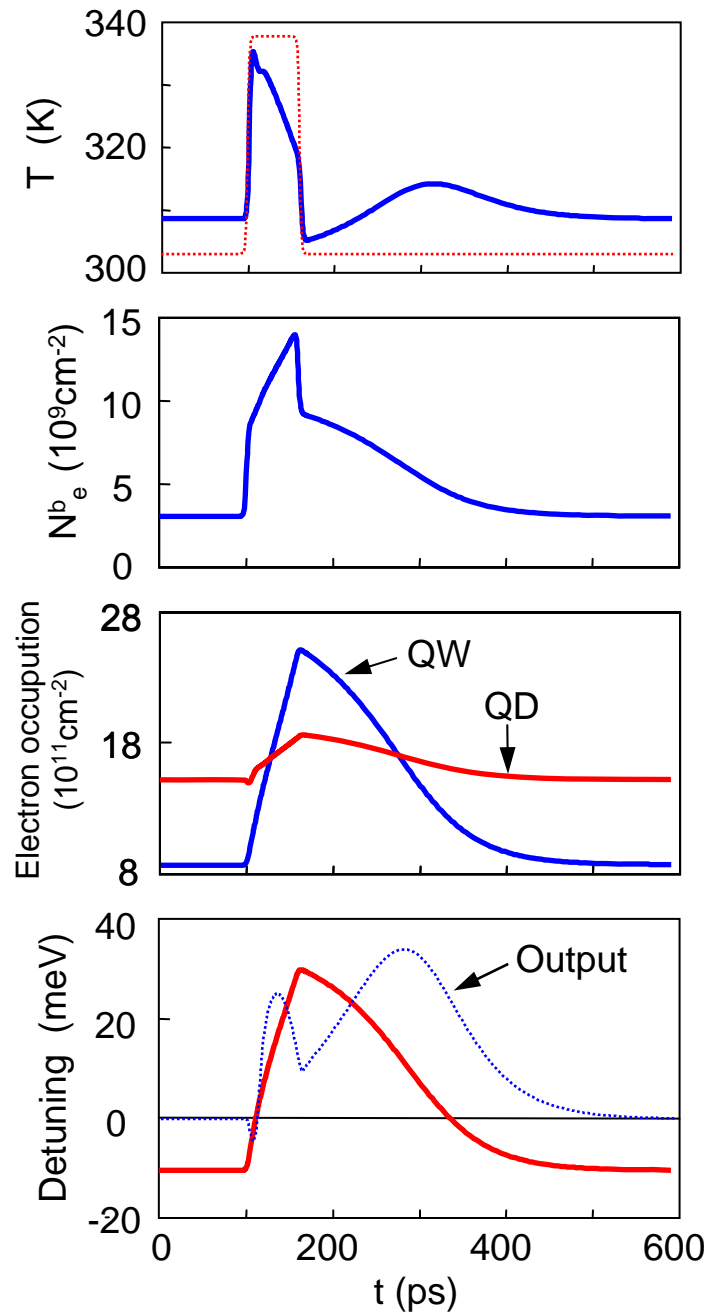
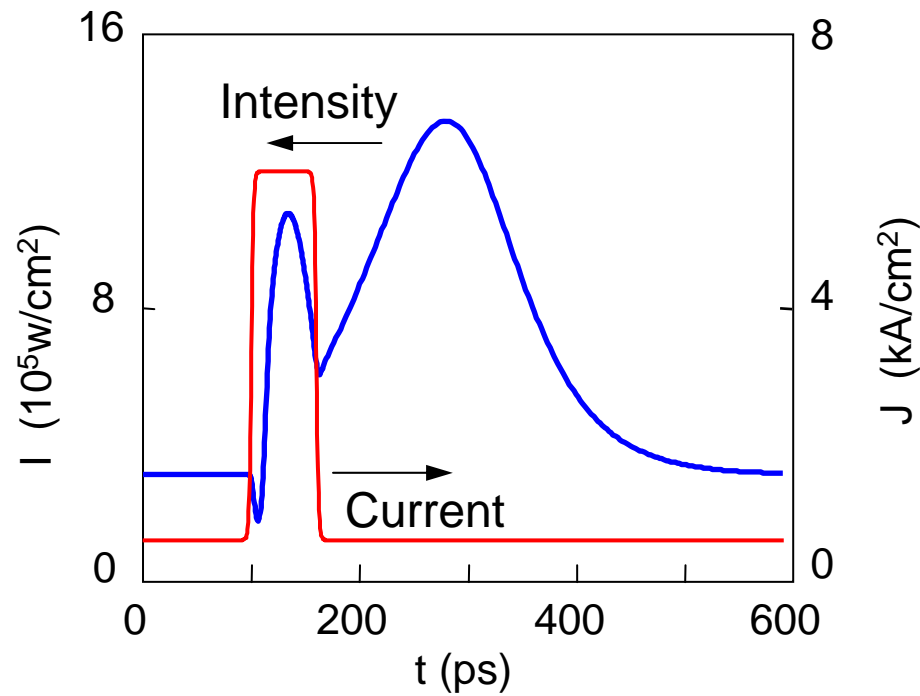
Microscopic theory in laser modeling



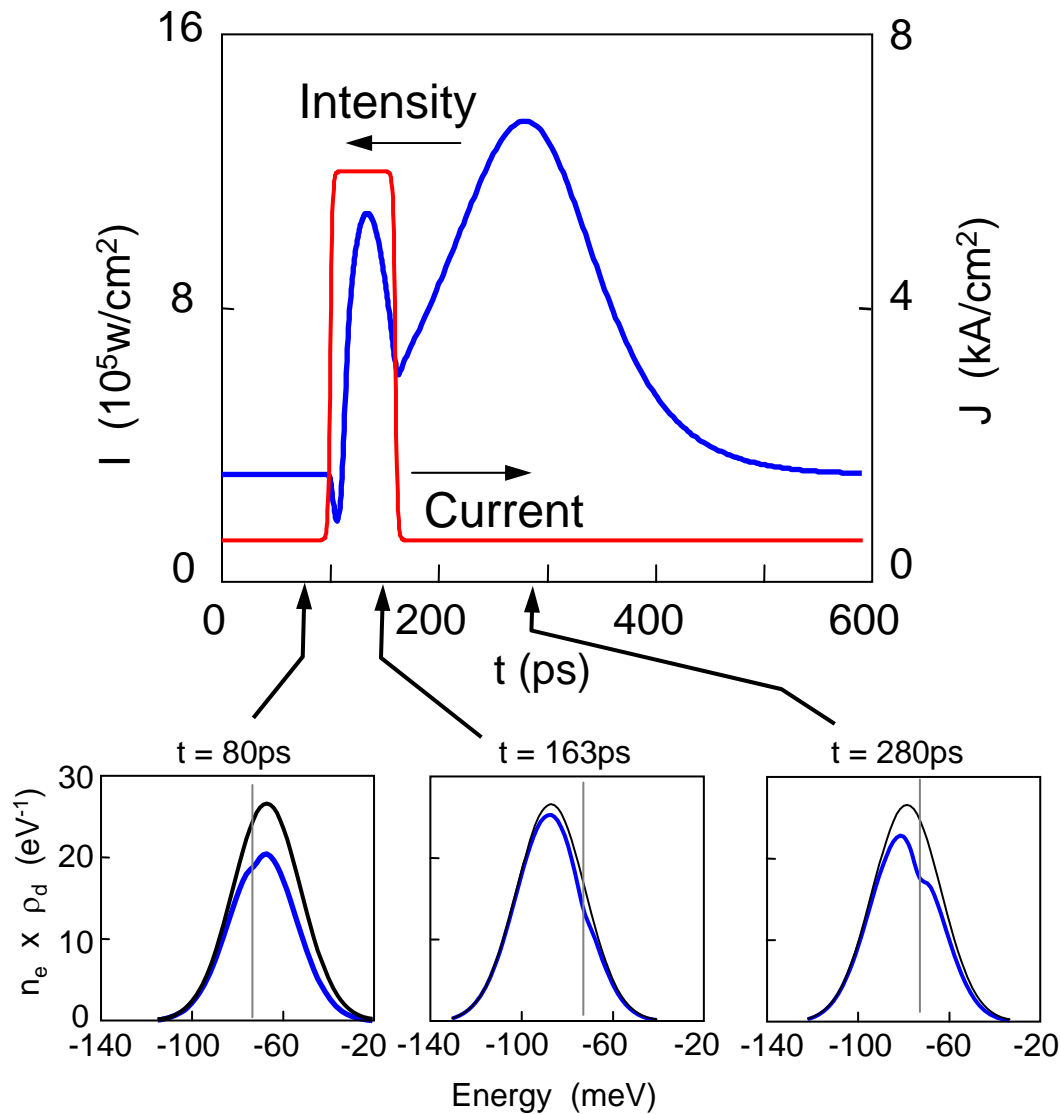
Nonequilibrium dynamics in a quantum-dot laser



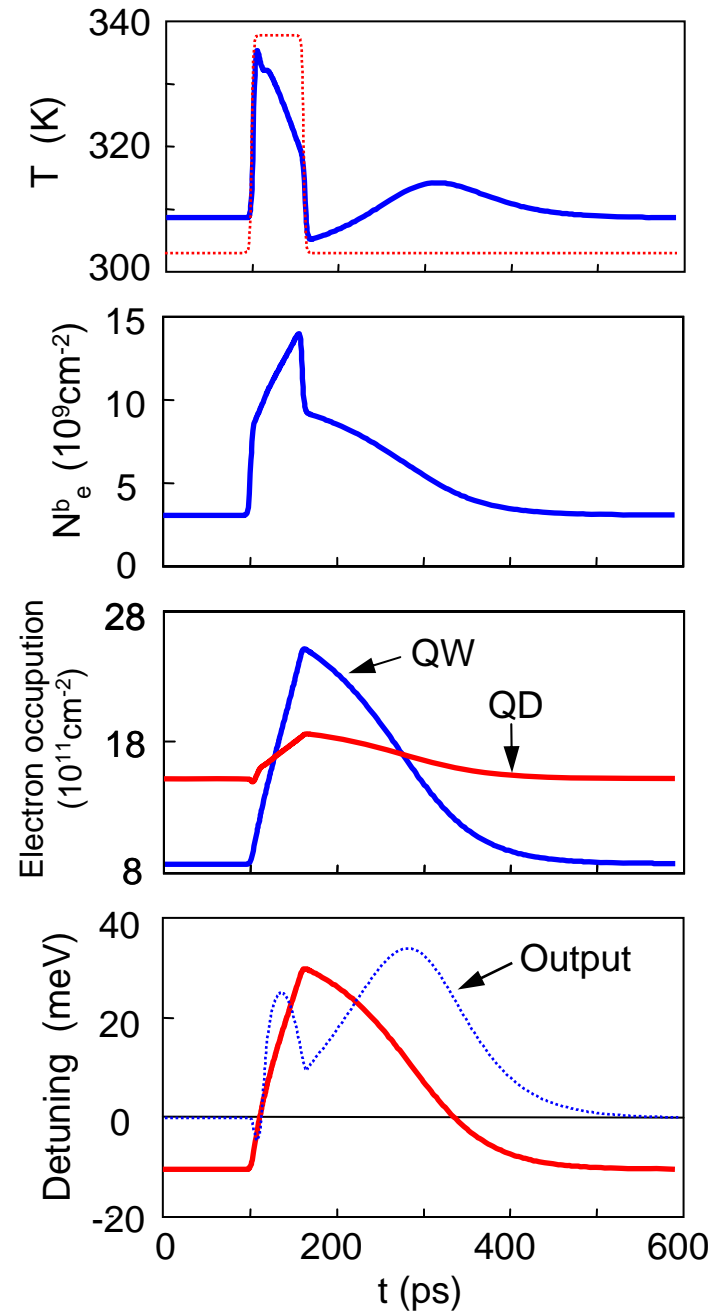
Nonequilibrium dynamics in a quantum-dot laser



Nonequilibrium dynamics in a quantum-dot laser



Hole burning in inhomogeneously-broadened quantum-dot population



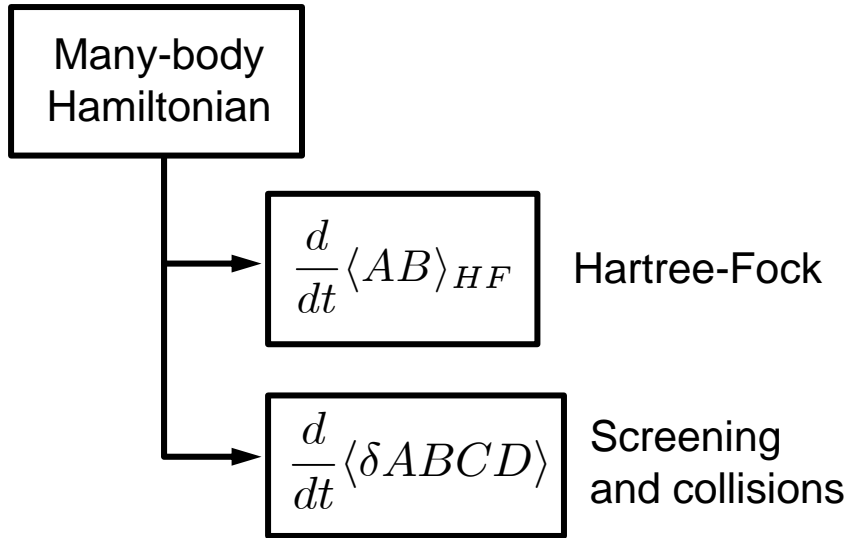
Summary

- 1) Past: Free-carrier theory
- 2) Present: Many-body description
- 3) Evolving: Quantum-dot gain theory

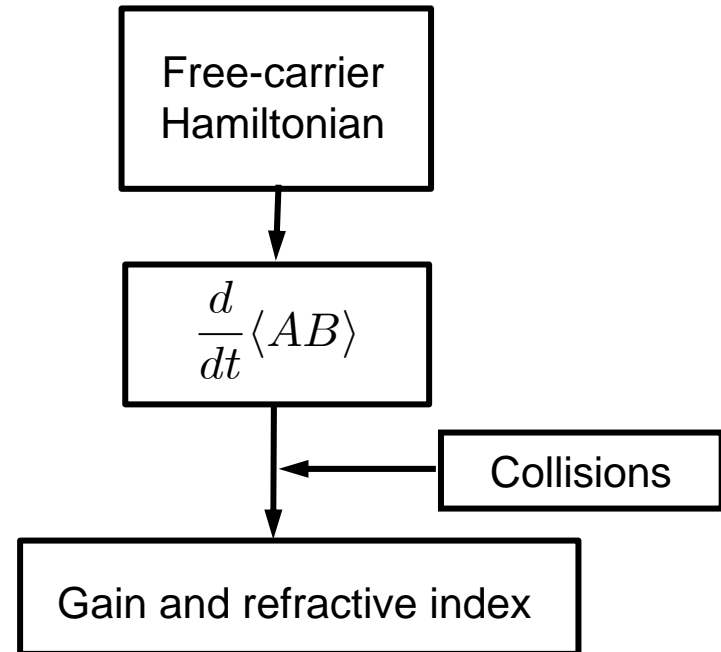
-
- Many-body effects \longleftrightarrow Coulomb interaction
 - Clean and understandable formulation
 - Influences all semiconductor gain properties
-

Schematic outline of gain approaches

Many body



Free carrier



Summary

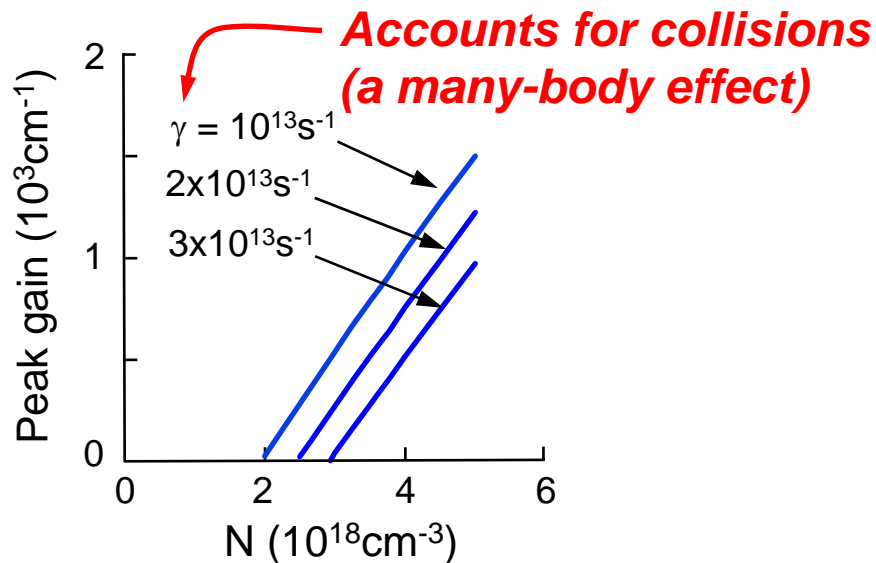
- 1) Past: Free-carrier theory
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- Many-body effects \longleftrightarrow Coulomb interaction
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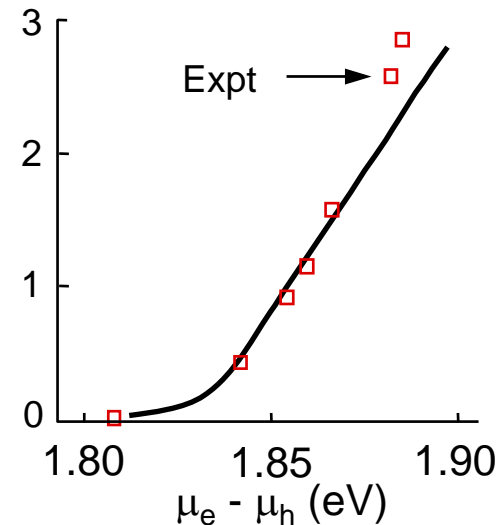
Free-carrier calculation



Many-body calculation



Peak gain sensitive to free parameter γ



Gain peak properties precisely and accurately determined

Uses of gain theory

Bandstructure
wurtzite or
zincblende,
bulk or
quantum well,
...

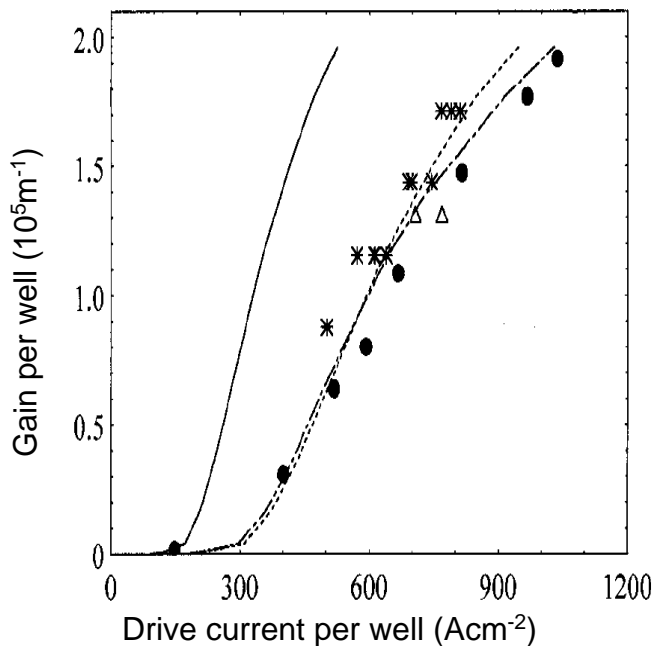
**Interacting
Plasma**
electrons
holes

**Gain and
refractive index**

**Laser
Field**

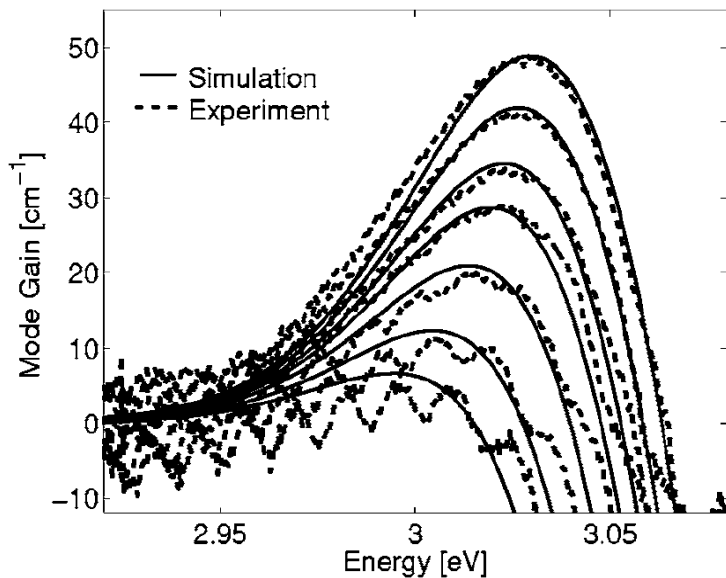
Extrinsic properties
inhomogeneous
broadening, carrier
leakage and capture,
nonradiative losses,

Expts



**Extraction of
extrinsic properties**

Wide bandgap Group-III Nitrides: very strong many-body effects

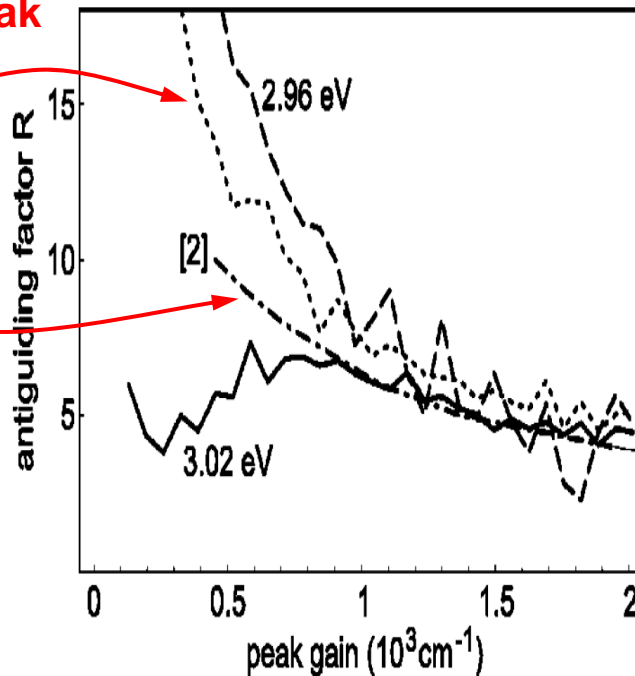


U.T. Schwarz (Univ. Regensburg), et al, APL 83, 4095 (2003)

α at gain peak

Expt.

Theory



B. Witzigmann (ETH Zürich), et al, APL 88, 021104 (2006)

$$g + iK_0\delta n = i \frac{\omega}{\epsilon_0 n c E} P$$

$$\alpha = R = - \frac{d(\delta n) / dN}{dg / dN}$$

α factor

Antiguiding factor

FIG. 4. Antiguiding factor R as function of peak gain at $E=3.02$ eV (solid curve), $E=2.96$ eV (dashed curve), and at the energy of the carrier-density-dependent peak gain (dotted curve). For comparison results from simulations for a single 2 nm QW is given (dot-dashed curve, from Ref. 2).

W.W.C, H. Amano and I. Akasaki, APL 76, 1647 (2000)

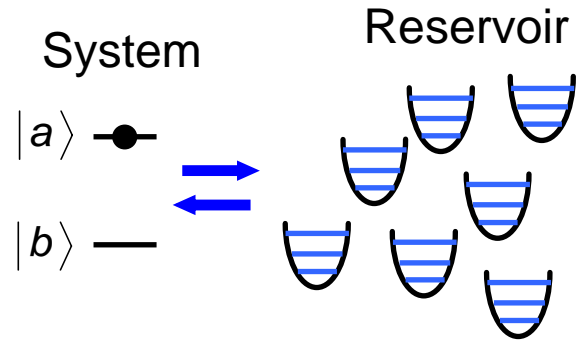
Polarization equation of motion

$$\underbrace{\omega_k^{(0)} - \sum_{k'} V_{kk'} (n_{ek'} + n_{hk'})}_{\text{Diagonal}} \quad \underbrace{\frac{2\mu_k E}{\hbar} + \sum_{k'} V_{kk'} p_{k'}}_{\text{Nondiagonal}}$$

$$\frac{dp_k}{dt} = -i\omega_k p_k - i\Omega_k (n_{ek} + n_{hk} - 1)$$

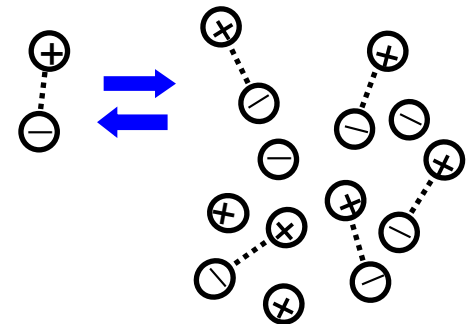
$$+ S_k^{c-c} + S_k^{c-p}$$

Typical



System Reservoir

Semiconductor



System is its own reservoir

$$- \sigma_k^d p_k + \sum_k \sigma_{k,k'}^{nd} p_{k'}$$

Diagonal

Nondiagonal

$$-p_\alpha \frac{2}{\hbar} \sum_{k'} \sum_{\mathbf{q}} \sum_{\sigma} \sum_{\sigma'} \sum_{\mathbf{k}} \xi_{\alpha k} \left[W_q^2 - \frac{\delta_{\sigma,\sigma'}}{2} W_q W_{|\mathbf{k}'-\mathbf{q}-\mathbf{k}|} \right] \frac{1}{\Delta - i\delta\varepsilon}$$

$$\times \left[n_{|\mathbf{k}'-\mathbf{q}|\sigma'} (1 - n_{k'\sigma'}) (1 - n_{k\sigma}) + (1 - n_{|\mathbf{k}'-\mathbf{q}|\sigma'}) n_{k'\sigma'} n_{k\sigma} \right]$$

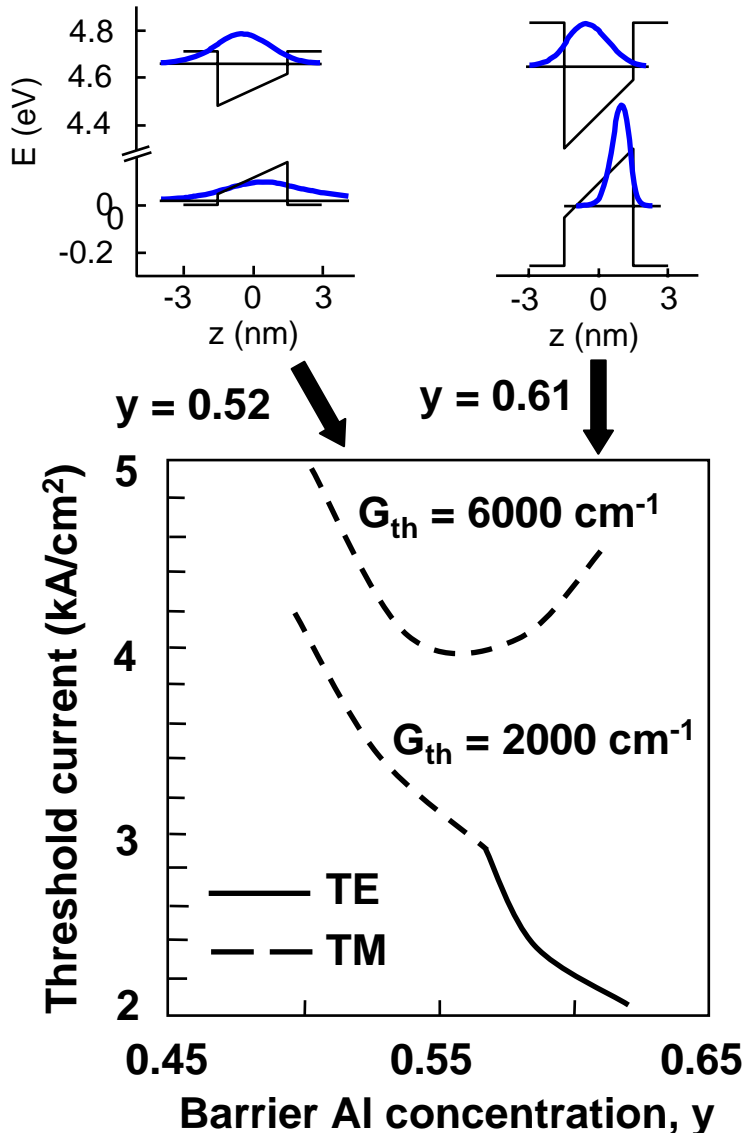
Barrier Influence on Threshold Current Density in AlGaNLasers

Chow (Sandia), Kneissl (Technical U, Berlin), Johnson and Northrup (PARC)

Weak QCSE

Strong QCSE

Simulation Model



- 3nm Al_{0.46}Ga_{0.54}N / Al_yGa_{1-y}N, 300K, 280 nm
- Many-body gain theory with carrier-carrier scattering treated at level of quantum kinetic theory (no T₂ free parameter)
- Band structure calculated with 6x6 **k·p** theory and Poisson equation
- Threshold current contributions: spontaneous emission, non-radiative losses and carrier leakage

Highlights

- Significant variation in threshold properties with barrier composition and threshold gain
- Gain structure optimization possible by changing only barrier configuration (and, hence, preserving emission wavelength)
- Predictive theory with only gain structure and bulk material parameters as input can facilitate laser design optimization

Computational Simulation of UV Nitride Laser Heterostructures (supporting information for simulation slide)

The graph in the lower left-hand quadrant illustrates the sensitivity of laser threshold current density to Al concentration, y , in an $\text{Al}_{0.46}\text{Ga}_{0.54}\text{N}$ quantum well structure with $\text{Al}_y\text{Ga}_{1-y}\text{N}$ barrier regions. The quantum well composition is chosen to give an emission wavelength of approximately 280 nm, which is useful for excitation of important biomolecules. The results suggest that optimization of laser gain configuration is possible by varying the barrier composition, with the quantum well width and composition remaining unchanged, thus preserving other desired laser properties such as the emission wavelength. The appreciable changes in threshold properties arise from changes in the quantum confinement Stark effect as illustrated in the top left-hand diagrams. The complicated behavior depicted in this slide is describable by a many-body theory that can be used to facilitate laser gain structure optimization for deep UV nitride laser heterostructures. Coordinating such computational optimization with growth of MQW heterostructures will accelerate progress toward the realization of deep UV laser diodes.